ME:5160 (58:160) Intermediate Mechanics of Fluids

Fall 2022 – HW9 Solution

P6.39 By analogy with laminar shear, $\tau = \mu du/dy$. T. V. Boussinesq in 1877 postulated that turbulent shear could also be related to the mean-velocity gradient $\tau \text{turb} = \varepsilon du/dy$, where ε is called the *eddy viscosity* and is much larger than μ . If the logarithmic-overlap law, Eq. (6.28), is valid with $\tau \approx \tau w$, show that $\varepsilon \approx \kappa \rho u^* y$.

Solution: Differentiate the log-law, Eq. (6.28), to find du/dy, then introduce the eddy viscosity into the turbulent stress relation

If
$$\frac{u}{u^*} = \frac{1}{\kappa} \ln\left(\frac{yu^*}{v}\right) + B$$
, then $\frac{du}{dy} = \frac{u^*}{\kappa y}$
Then, if $\tau \approx \tau_w = \rho u^{*2} = \varepsilon \frac{du}{dy} = \varepsilon \frac{u^*}{\kappa y}$, solve for $\varepsilon = \kappa \rho u^* y$ Ans.

Note that $\varepsilon/\mu = \kappa y^+$, which is much larger than unity in the overlap region.

P6.40 Theodore von Kármán in 1930 theorized that turbulent shear could be represented by τ turb $= \varepsilon du/dy$ where $\varepsilon = \rho \kappa^2 y^2 |du/dy|$ is called the *mixing-length eddy viscosity* and $\kappa \approx 0.41$ is Kármán's dimensionless *mixing-length constant* [2,3]. Assuming that τ turb $\approx \tau w$ near the wall, show that this expression can be integrated to yield the logarithmic-overlap law, Eq. (6.28).

Solution: This is accomplished by straight substitution:

$$\tau_{\text{turb}} \approx \tau_{\text{w}} = \rho u^{*2} = \varepsilon \frac{\text{d}u}{\text{d}y} = \left[\rho \kappa^2 y^2 \left| \frac{\text{d}u}{\text{d}y} \right| \right] \frac{\text{d}u}{\text{d}y}, \text{ solve for } \frac{\text{d}u}{\text{d}y} = \frac{u^*}{\kappa y}$$

Integrate:
$$\int du = \frac{u^*}{\kappa} \int \frac{dy}{y}$$
, or: $u = \frac{u^*}{\kappa} \ln(y) + \text{constant}$ Ans.

To convert this to the exact form of Eq. (6.28) requires fitting to experimental data

P6.45 Oil, SG = 0.88 and $v = 4E-5 \text{ m}^2/\text{s}$, flows at 400 gal/min through a 6-inch asphalted castiron pipe. The pipe is 0.5 miles long (2640 ft) and slopes upward at 8° in the flow direction. Compute the head loss in feet and the pressure change.

Solution: First convert 400 gal/min = 0.891 ft³/s and $\nu = 0.000431$ ft²/s. For asphalted cast-iron, $\varepsilon = 0.0004$ ft, hence $\varepsilon/d = 0.0004/0.5 = 0.0008$. Compute *V*, *Red*, and *f*:

$$V = \frac{0.891}{\pi (0.25)^2} = 4.54 \ \frac{ft}{s}; \quad Re_d = \frac{4.54(0.5)}{0.000431} = 5271; \quad \text{calculate} \quad f_{Moody} = 0.0377$$

then $h_f = f \frac{L}{d} \frac{V^2}{2g} = 0.0377 \left(\frac{2640}{0.5}\right) \frac{(4.54)^2}{2(32.2)} = 63.8 \text{ ft} \quad Ans. \text{ (a)}$

If the pipe slopes upward at 8°, the pressure drop must balance both friction and gravity:

$$\Delta p = \rho g(h_f + \Delta z) = 0.88(62.4)[63.8 + 2640\sin 8^\circ] = 23700 \frac{101}{\text{ft}^2} \quad Ans. \text{ (b)}$$

P6.61 What level *h* must be maintained in Fig. P6.61 to deliver a flow rate of 0.015 ft³/s through the $\frac{1}{2}$ -in commercial-steel pipe?



Solution: For water at 20°C, take $\rho = 1.94$ slug/ft³ and $\mu = 2.09E-5$ slug/ft·s. For commercial steel, take $\varepsilon \approx 0.00015$ ft, or $\varepsilon/d = 0.00015/(0.5/12) \approx 0.0036$. Compute

$$V = \frac{Q}{A} = \frac{0.015}{(\pi/4)(0.5/12)^2} = 11.0 \frac{ft}{s};$$

$$\operatorname{Re} = \frac{\rho \operatorname{Vd}}{\mu} = \frac{1.94(11.0)(0.5/12)}{2.09E-5} \approx 42500 \quad \varepsilon/d = 0.0036, \quad \operatorname{f}_{\operatorname{Moody}} \approx 0.0301$$

The energy equation, with $p_1 = p_2$ and $V_1 \approx 0$, yields an expression for surface elevation:

$$h = h_{f} + \frac{V^{2}}{2g} = \frac{V^{2}}{2g} \left(1 + f\frac{L}{d}\right) = \frac{(11.0)^{2}}{2(32.2)} \left[1 + 0.0301 \left(\frac{80}{0.5/12}\right)\right] \approx 111 \text{ ft} \quad Ans.$$