# ME:5160 (58:160) Intermediate Mechanics of Fluids <br> Fall 2022 - HW13 Solution 

P8.15 Hurricane Sandy, which hit the New Jersey coast on Oct. 29, 2012, was extremely broad, with wind velocities of $40 \mathrm{mi} / \mathrm{h}$ at 400 miles from its center. Its maximum velocity was $90 \mathrm{mi} / \mathrm{h}$. Using the model of Fig. P8.14, at $20^{\circ} \mathrm{C}$ with a pressure of 100 kPa far from the center, estimate (a) the radius $R$ of maximum velocity, in mi; and (b) the pressure at $r=R$.

Solution: The air density is $p / R T=(100,000) /[287(293)]=1.19 \mathrm{~kg} / \mathrm{m}^{3}$. Convert $90 \mathrm{mi} / \mathrm{h}$ to $40.23 \mathrm{~m} / \mathrm{s}$ and $40 \mathrm{mi} / \mathrm{h}$ to $17.9 \mathrm{~m} / \mathrm{s}$. The outer flow is irrotational, hence Bernoulli holds:

$$
\begin{aligned}
p_{r=R}+\frac{1}{2} \rho V_{r=R}^{2}=p_{r=R}+\frac{1}{2}(1.19)(40.2)^{2}=p_{\infty}+\frac{1}{2} \rho V_{\infty}^{2} & =100,000+0 \\
\text { solve for } \quad p_{r=R}=100,000-960=99,040 P a & \approx \mathbf{9 9 k} \mathbf{k P a} \quad \text { Ans. }(b)
\end{aligned}
$$

(a) The two velocities, plus the radii, fit irrotational vortex theory. Let the outer ring be $r_{2}$ :

$$
V_{2}=40 m i / h=\frac{C}{r_{2}}=\frac{C}{400 m i} \text {, hence } C=(40)(400)=16,000 \mathrm{mi}^{2} / h
$$

Then $\quad V_{R}=90 \mathrm{mi} / \mathrm{h}=\frac{16,000 \mathrm{~m}^{2} / \mathrm{h}}{R}$, solve $R \approx \mathbf{1 7 8} \mathbf{m i}$ Ans.(a)

P8.27 Water at $20^{\circ} \mathrm{C}$ flows past a half-body as shown in Fig. P8.27. Measured pressures at points A and B are 160 kPa and 90 kPa , respectively, with uncertainties of 3 kPa each. Estimate the stream velocity and its uncertainty.

Solution: Since Eq. (8.18) is for the upper surface, use it by noting that $V \mathrm{C}=V \mathrm{~B}$ in the figure:


Fig. P8. 27

$$
\begin{aligned}
& \frac{r_{C}}{a}=\frac{\pi-\pi / 2}{\sin (\pi / 2)}=\frac{\pi}{2}, \quad V_{C}^{2}=V_{B}^{2}=U_{\infty}^{2}\left[1+\left(\frac{2}{\pi}\right)^{2}+\frac{2}{(\pi / 2)} \cos (\pi / 2)\right]=1.405 U_{\infty}^{2} \\
& \text { Bernoulli: } \quad p_{A}+\frac{\rho}{2} V_{A}^{2}=160000+0=p_{B}+\frac{\rho}{2} V_{B}^{2}=90000+\frac{998}{2}\left(1.405 U_{\infty}^{2}\right)
\end{aligned}
$$

Solve for $\mathrm{U}_{\infty} \approx \mathbf{1 0 . 0} \mathbf{m} / \mathrm{s}$ Ans.
The uncertainty in $\left(p_{\mathrm{A}}-p_{\mathrm{B}}\right)$ is as high as 6000 Pa , hence the uncertainty in $U_{\infty}$ is $\pm \mathbf{0 . 4} \mathbf{~ m} / \mathrm{s}$. Ans.

P8.29 A uniform water stream, $U_{\infty}=20 \mathrm{~m} / \mathrm{s}$ and $\rho=998 \mathrm{~kg} / \mathrm{m}^{3}$, combines with a source at the origin to form a half-body. At $(x, y)=(0,1.2 \mathrm{~m})$, the pressure is 12.5 kPa less than $p_{\infty}$. (a) Is this point outside the body? Estimate (b) the appropriate source strength $m$ and (c) the pressure at the nose of the body.


Fig P8.29
Solution: We know, from Fig. 8.5 and Eq. 8.18, the point on the half-body surface just above " $m$ " is at $\mathrm{y}=\pi \mathrm{a} / 2$, as shown, where $\mathrm{a}=m / \mathrm{U}$. The Bernoulli equation allows us to compute the necessary source strength $m$ from the pressure at $(x, y)=(0,1.2 \mathrm{~m})$ :

$$
\mathrm{p}_{\infty}+\frac{\rho}{2} \mathrm{U}_{\infty}^{2}=\mathrm{p}_{\infty}+\frac{998}{2}(20)^{2}=\mathrm{p}_{\infty}-12500+\frac{998}{2}\left[(20)^{2}+\left(\frac{m}{1.2}\right)^{2}\right]
$$

Solve for $\boldsymbol{m}-6.0 \frac{\mathbf{m}^{2}}{\mathbf{s}} \quad$ Ans. (b) while $\mathrm{a}=\frac{m}{\mathrm{U}}=\frac{6.0}{20}=0.3 \mathrm{~m}$

The body surface is thus at $\mathrm{y}=\pi \mathrm{a} / 2=\mathbf{0 . 4 7} \mathbf{~ m}$ above m . Thus the point in question, $\mathrm{y}=1.2 \mathrm{~m}$ above $m$, is outside the body. Ans. (a)

At the nose SP of the body, $(x, y)=(-a, 0)$, the velocity is zero, hence we predict

$$
\mathrm{p}_{\infty}+\frac{\rho}{2} \mathrm{U}_{\infty}^{2}=\mathrm{p}_{\infty}+\frac{998}{2}(20)^{2}=\mathrm{p}_{\text {nose }}+\frac{\rho}{2}(0)^{2}, \quad \text { or } \quad \mathbf{p}_{\text {nose }} \approx \mathbf{p}_{\infty}+\mathbf{2 0 0} \mathbf{k P a} \quad \text { Ans. (c) }
$$

P8.44 Suppose that circulation is added to the cylinder flow of Prob. P8.43 sufficient to place the stagnation points at $\theta=35^{\circ}$ and $145^{\circ}$. What is the required vortex strength $K$ in $\mathrm{m}^{2} / \mathrm{s}$ ? Compute the resulting pressure and surface velocity at (a) the stagnation points, and (b) the upper and lower shoulders. What will be the lift per meter of cylinder width?

Solution: Recall that Prob. P8.43 was for water at $20^{\circ} \mathrm{C}$ flowing at $6 \mathrm{~m} / \mathrm{s}$ past a $1-\mathrm{m}$-diameter cylinder, with $p_{\infty}=200 \mathrm{kPa}$. From Eq. (8.35),

$$
\sin \theta_{\text {stag }}=\sin \left(35^{\circ}\right)=\frac{K}{2 U_{\infty} a}=\frac{K}{2(6 \mathrm{~m} / \mathrm{s})(0.5 \mathrm{~m})}, \quad \text { or: } \quad K=3.44 \mathrm{~m}^{2} / \mathrm{s} \quad \text { Ans. }
$$

(a) At the stagnation points, velocity is zero and pressure equals stagnation pressure:

$$
p_{\text {stag }}=p_{\infty}+\frac{\rho}{2} U_{\infty}^{2}=200,000 P a+\frac{998 \mathrm{~kg} / \mathrm{m}^{3}}{2}(6 \mathrm{~m} / \mathrm{s})^{2}=\mathbf{2 1 8 , 0 0 0} \mathbf{P a} \quad \text { Ans. (a) }
$$

(b) At any point on the surface, from Eq. (8.37),

$$
p_{\text {stag }}=218000=p_{\text {surf }}+\frac{\rho}{2}\left(-2 U_{\infty} \sin \theta+\frac{K}{a}\right)^{2}=p_{\text {surf }}+\frac{998}{2}\left[-2(6) \sin \theta+\frac{3.44}{0.5}\right]^{2}
$$

At the upper shoulder, $\theta=90^{\circ}$,

$$
p=218000-\frac{998}{2}(-5.12)^{2} \approx \mathbf{2 0 4 , 9 0 0} \mathbf{P a} \text { Ans. (b-upper) }
$$

At the lower shoulder, $\theta=270^{\circ}$,

$$
p=218000-\frac{998}{2}(-18.88)^{2} \approx \mathbf{4 0 , 1 0 0} \mathbf{P a} \text { Ans. (b-lower) }
$$

*P8.48 Wind at $\mathrm{U}_{\infty}$ and $\mathrm{p}_{\infty}$ flows past a Quonset hut which is a half-cylinder of radius a and length L (Fig. P8.48). The internal pressure is pi. Using inviscid theory, derive an expression for the upward force on the hut due to the difference between pi and ps.


Fig P8.48
Solution: The analysis is similar to Prob. P8.46 on the previous page. If po is the stagnation pressure at the nose $\left(\theta=180^{\circ}\right)$, the surface pressure distribution is

$$
\mathrm{p}_{\mathrm{s}}=\mathrm{p}_{\mathrm{o}}-\frac{\rho}{2} \mathrm{U}_{\mathrm{s}}^{2}=\mathrm{p}_{\mathrm{o}}-\frac{\rho}{2}\left(2 \mathrm{U}_{\infty} \sin \theta\right)^{2}=\mathrm{p}_{\mathrm{o}}-2 \rho \mathrm{U}_{\infty}^{2} \sin ^{2} \theta
$$

Then the net upward force on the half-cylinder is found by integration:

$$
\begin{aligned}
& \mathrm{F}_{\mathrm{up}}=\int_{0}^{\pi}\left(\mathrm{p}_{\mathrm{i}}-\mathrm{p}_{\mathrm{s}}\right) \sin \theta \mathrm{abd} \theta=\int_{0}^{\pi}\left(\mathrm{p}_{\mathrm{i}}-\mathrm{p}_{\mathrm{o}}+2 \rho \mathrm{U}_{\infty}^{2} \sin ^{2} \theta\right) \sin \theta \mathrm{abd} \theta, \\
& \text { or: } \quad \mathbf{F}_{\mathrm{up}}=\left(\mathbf{p}_{\mathbf{i}}-\mathbf{p}_{\mathbf{o}}\right) \mathbf{2 a b}+\frac{\mathbf{8}}{\mathbf{3}} \rho \mathbf{U}_{\infty}^{\mathbf{2}} \mathbf{a b} \quad \text { Ans. }\left(\text { where } \mathrm{p}_{\mathrm{o}}=\mathrm{p}_{\infty}+\frac{\rho}{2} \mathrm{U}_{\infty}^{2}\right)
\end{aligned}
$$

P8.75 Using the four-source image pattern needed to construct the flow near a corner shown in Fig. P8.72, find the value of the source strength $m$ which will induce a wall velocity of $4.0 \mathrm{~m} / \mathrm{s}$ at the point $(x, y)=(a, 0)$ just below the source shown, if $a=50 \mathrm{~cm}$.

Solution: The flow pattern is formed by four equal sources $m$ in the 4 quadrants, as in the figure at right. The sources above and below the point $\mathrm{A}(a, 0)$ cancel each other at A , so the velocity at A is caused only by the two left sources. The velocity at A is the sum of the two horizontal components from these 2 sources:


Fig P8.75
$V_{A}=2 \frac{m}{\sqrt{a^{2}+(2 a)^{2}}} \frac{2 a}{\sqrt{a^{2}+(2 a)^{2}}}=\frac{4 m a}{5 a^{2}}=\frac{4 m}{5(0.5 m)}=4 \frac{\mathrm{~m}}{\mathrm{~s}} \quad$ if $\boldsymbol{m}=\mathbf{2 . 5} \frac{\mathbf{m}^{2}}{\mathbf{s}} \quad$ Ans.

C8.4 Find a formula for the stream function for flow of a doublet of strength $\lambda$ at a distance $a$ from a wall, as in Fig. C8.4. (a) Sketch the streamlines. (b) Are there any stagnation points? (c) Find the maximum velocity along the wall and its position.


Fig C8.4
Solution: Use an image doublet of the same strength and orientation at the $(x, y)=(0,-a)$. The stream function for this combined flow will form a "wall" at $y=0$ between the two doublets:

$$
\psi=-\frac{\lambda(y+a)}{x^{2}+(y+a)^{2}}-\frac{\lambda(y-a)}{x^{2}+(y-a)^{2}}
$$

(a) The streamlines are shown on the next page for one quadrant of the doubly-symmetric flow field. They are fairly circular, like Fig. 8.8, above the doublet, but they flatten near the wall.


## Problem C8.4

(b) There are no stagnation points in this flow field. Ans. (b)
(c) The velocity along the wall $(y=0)$ is found by differentiating the stream function:

$$
u_{\text {wall }}=\left.\frac{\partial \psi}{\partial y}\right|_{y=0}=-\frac{\lambda}{x^{2}+a^{2}}+\frac{2 \lambda a^{2}}{\left(x^{2}+a^{2}\right)^{2}}-\frac{\lambda}{x^{2}+a^{2}}+\frac{2 \lambda a^{2}}{\left(x^{2}+a^{2}\right)^{2}}
$$

The maximum velocity occurs at $x=0$, that is, right between the two doublets:

$$
u_{w, \max }=\frac{2 \lambda}{a^{2}} \quad \text { Ans. (c) }
$$

