

Review for Exam3

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Chapter 8

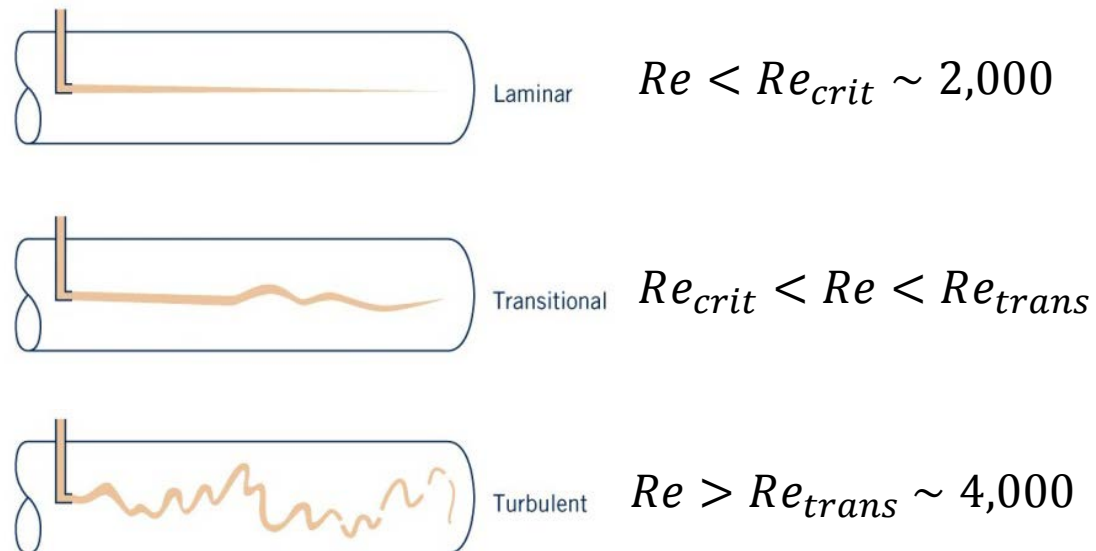
Flow in Conduits

- Internal flow: Confined by solid walls
- Basic piping problems:
 - Given the desired flow rate, what pressure drop is needed to drive the flow?
 - Given the pressure drop available, what flow rate will ensue?

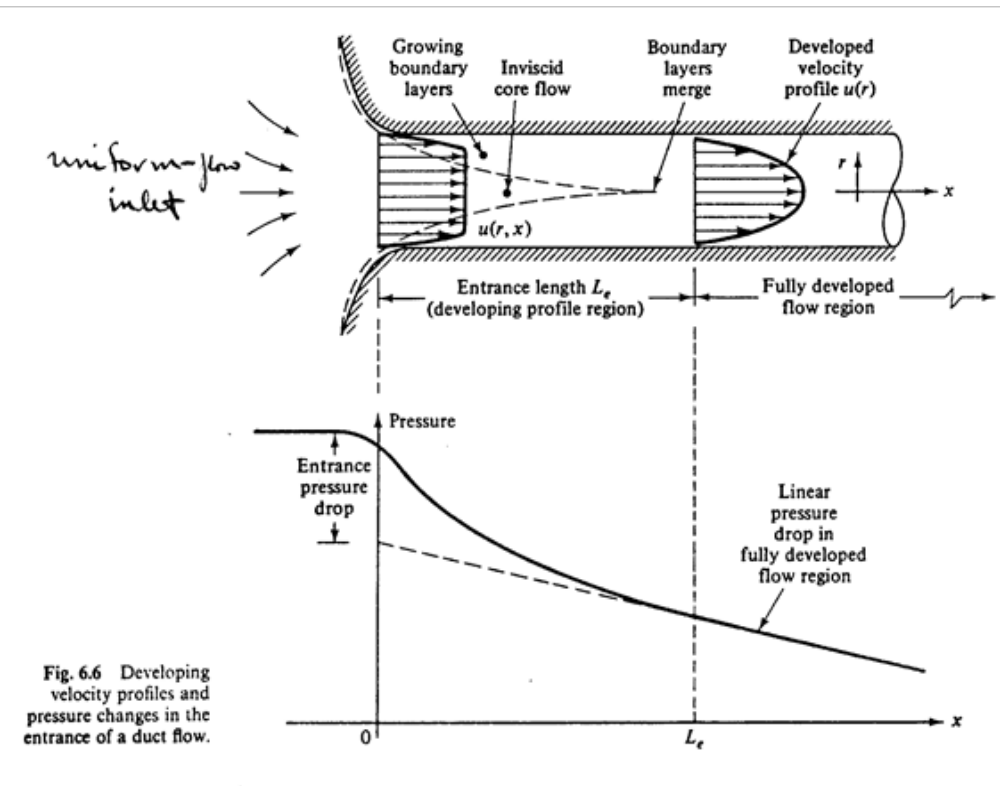
Pipe Flow: Laminar vs. Turbulent

- Reynolds number regimes

$$Re = \frac{\rho V D}{\mu}$$



Entrance Region and Fully Developed



- Entrance Length, L_e :

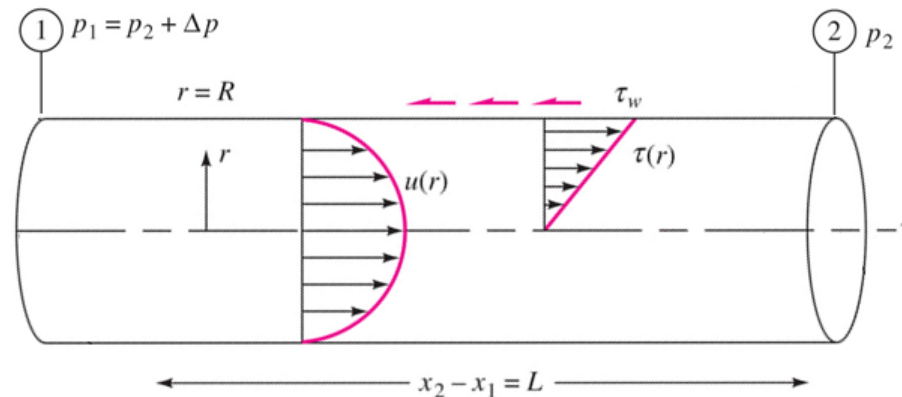
- Laminar flow: $Le/D = 0.06Re$ ($L_{e,max} = 0.06Re_{crit} \sim 138D$)

- Turbulent flow: $Le/D = 4.4Re^{\frac{1}{6}}$ ($20D < L_e < 30D$ for $10^4 < Re < 10^5$)

Pressure Drop and Shear Stress

- Pressure drop, $\Delta p = p_1 - p_2$, is needed to overcome viscous shear stress.
- Considering force balance,

$$p_1 \left(\frac{\pi D^2}{4} \right) - p_2 \left(\frac{\pi D^2}{4} \right) = \tau_w (\pi D L) \Rightarrow \Delta p = 4 \tau_w \frac{L}{D}$$



Head Loss and Friction Factor

- Energy equation

$$h_L = \frac{p_1 - p_2}{\gamma} + \frac{\alpha_1 V_1^2 - \alpha_2 V_2^2}{2g} + (z_1 - z_2) = \frac{\Delta p}{\gamma}$$

$$\Delta p = 4\tau_w \frac{L}{D} \quad \text{from force balance and } \gamma = \rho g$$

$$\therefore h_L (\text{or } h_f) = 4\tau_w \frac{L}{D} / \rho g = \left(\frac{8\tau_w}{\rho V^2} \right) \cdot \frac{L V^2}{D 2g} = f \frac{L V^2}{D 2g}$$

\Rightarrow Darcy – Weisbach equation

- Friction factor

$$f \equiv \frac{8\tau_w}{\rho V^2}$$

Fully-developed Laminar Flow

- Exact solution, $u(r) = V_c \left[1 - \left(\frac{r}{R} \right)^2 \right]$

- Wall shear stress

$$\tau_w = -\mu \left. \frac{du}{dr} \right|_{r=R} = \frac{8\mu V}{D}$$

Where, $V = Q/A$

- Friction factor,

$$f = \frac{8\tau_w}{\rho V^2} = \frac{8}{\rho V^2} \cdot \frac{8\mu V}{D} = \frac{64}{\rho D V / \mu} = \frac{64}{Re}$$

Fully-developed Turbulent Flow

- Dimensional analysis

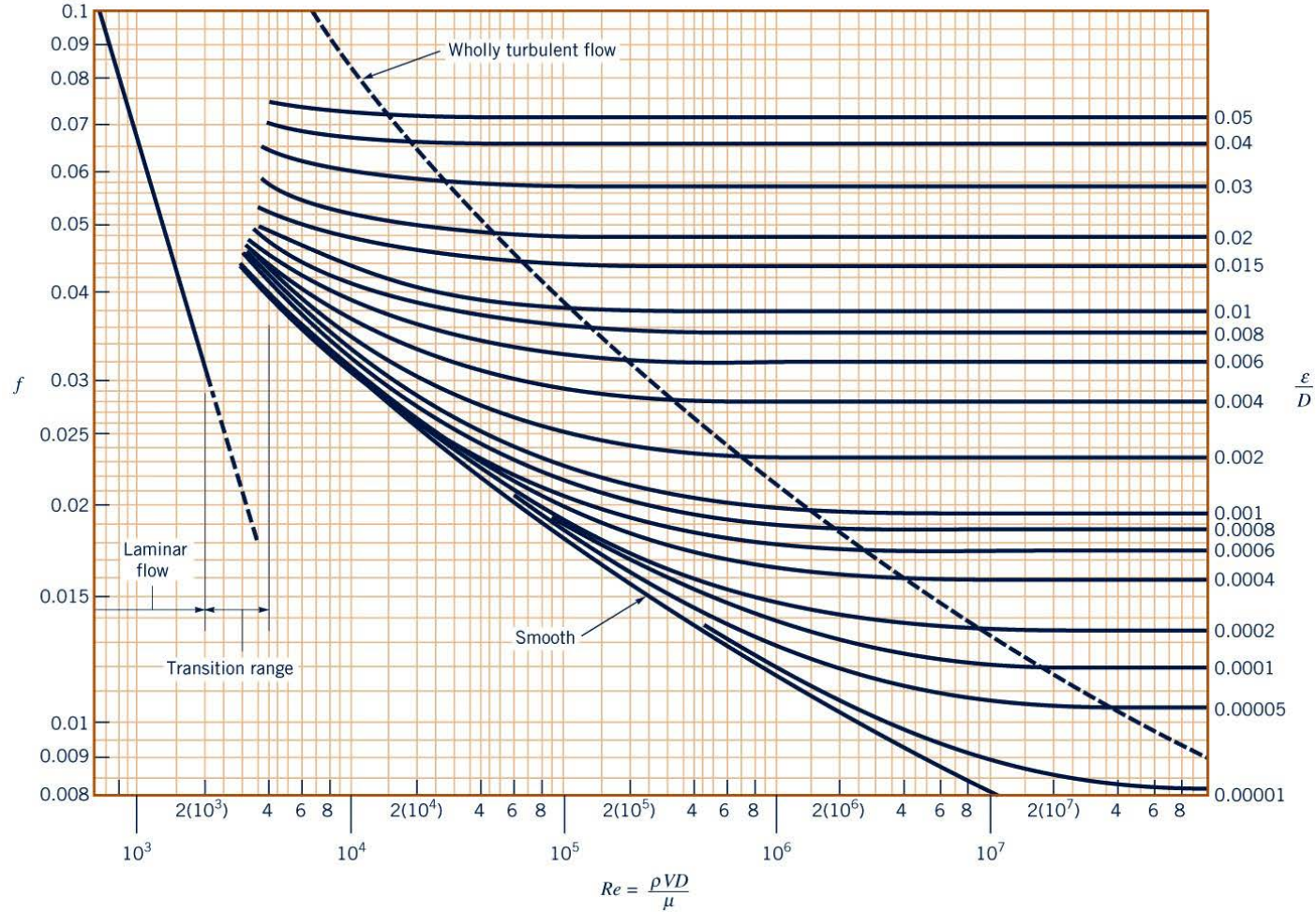
$$\tau_w = f(D, V, \mu, \rho, \varepsilon)$$

$$\rightarrow k - r = 6 - 3 = 3 \Pi's$$

$$\frac{\tau_w}{\rho V^2} = \phi\left(\frac{\rho V D}{\mu}, \frac{\varepsilon}{D}\right)$$

$$\therefore f = \phi(Re, \varepsilon/D)$$

Moody Chart



Moody Chart – Contd.

- Colebrook equation

$$\frac{1}{\sqrt{f}} = -2 \log \left(\frac{\varepsilon/D}{3.7} + \frac{2.51}{Re\sqrt{f}} \right)$$

- Haaland equation

$$\frac{1}{\sqrt{f}} = -1.8 \log \left[\left(\frac{\varepsilon/D}{3.7} \right)^{1.1} + \frac{6.9}{Re} \right]$$

Minor Loss

- Loss of energy due to pipe system components (valves, bends, tees, and the like).
- Theoretical prediction is, as yet, almost impossible.
- Usually based on experimental data.

K_L : Loss coefficient

$$h_m = \sum K_L \frac{V^2}{2g}$$

Pipe Flow Problems

- Energy equation for pipe flow:

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + z_1 + h_p = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + z_2 + h_t + \left(f \frac{L}{D} + \sum K_L \right) \frac{V^2}{2g}$$

- Type I: Determine head loss h_L (or Δp)
- Type II: Determine flow rate Q (or V)
- Type III: Determine pipe diameter D

Iteration is needed for types II and III

Type I Problem

- Typically, V (or Q) and D are given $\rightarrow Re$ and ε/D

$$f = \phi \left(\frac{\rho V D}{\mu}, \frac{\varepsilon}{D} \right) \Rightarrow \text{from Moody Chart}$$

$$\therefore h_L = f \frac{L}{D} \frac{V^2}{2g}$$

Type II Problem

- Q (thus V) is unknown $\rightarrow Re?$
- Solve energy equation for V as a function of f . For example, if $p_1 = p_2$, $V_1 = V_2$ and $\Delta z = z_1 - z_2$,

$$0 + 0 + (z_1 - z_2) + h_p = 0 + 0 + \left(f \frac{L}{D} + \sum K_L \right) \frac{V^2}{2g}$$

$$\therefore V = \sqrt{\frac{2g(h_p + \Delta z)}{f \frac{L}{D} + \sum K_L}}$$

Guess $f \rightarrow V \rightarrow Re \rightarrow f_{\text{new}}$; Repeat until f is converged $\Rightarrow V$

Type III Problem

- D is unknown $\rightarrow Re$ and ε/D ?
- Solve energy equation for D as a function of f . For example, if $p_1 = p_2$, $V_1 = V_2$, $\Delta z = z_1 - z_2$, and $\sum K_L = 0$ and using $V = Q/(\pi D^2/4)$,

$$0 + 0 + (z_1 - z_2) + h_p = 0 + 0 + \left(f \frac{L}{D} + 0\right) \frac{1}{2g} \left(\frac{Q}{\pi D^2/4}\right)^2$$

$$\therefore D = \left[\frac{8LQ^2 \cdot f}{\pi^2 g (h_p + \Delta z)} \right]^{\frac{1}{5}}$$

Guess $f \rightarrow D \rightarrow Re$ and $\varepsilon/D \rightarrow f_{\text{new}}$; Repeat until f is converged $\Rightarrow D$

Chapter 9

Flow over Immersed Bodies

- External flow: Unconfined, free to expand
- Complex body geometries require experimental data (dimensional analysis)

Drag

- Resultant force in the direction of the upstream velocity

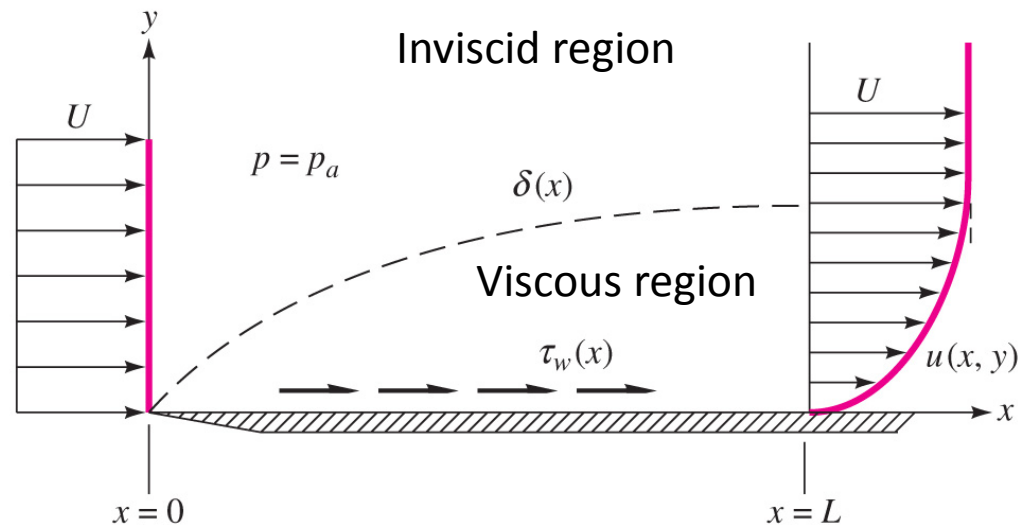
$$C_D = \frac{D}{\frac{1}{2}\rho V^2 A} = \frac{1}{\frac{1}{2}\rho V^2 A} \left(\underbrace{\int_S (p - p_\infty) \underline{n} \cdot \hat{\mathbf{i}} dA}_{C_{Dp} = \text{Pressure drag (or Form drag)}} + \underbrace{\int_S \tau_w \underline{t} \cdot \hat{\mathbf{i}} dA}_{C_f = \text{Friction drag}} \right)$$

- Streamlined body ($t/\ell \ll 1$): $C_f \gg C_{Dp}$, Boundary layer flow
- Bluff body ($t/\ell \sim 1$): $C_{Dp} \gg C_f$

where, t is the thickness and ℓ the length of the body

Boundary Layer

- High Reynolds number flow, $Re_x = \frac{Ux}{\nu} \gg 1,000$
- Viscous effects are confined to a thin layer, δ
- $\frac{u}{U} = 0.99$ at $y = \delta$



Friction Coefficient

- Local friction coefficient

$$c_f(x) = \frac{2\tau_w(x)}{\rho U^2}$$

Note: Darcy friction factor for pipe flow

$$f = \frac{8\tau_w}{\rho V^2}$$

- Friction drag coefficient

$$D_f = \int_{A=bL} \tau_w dA$$

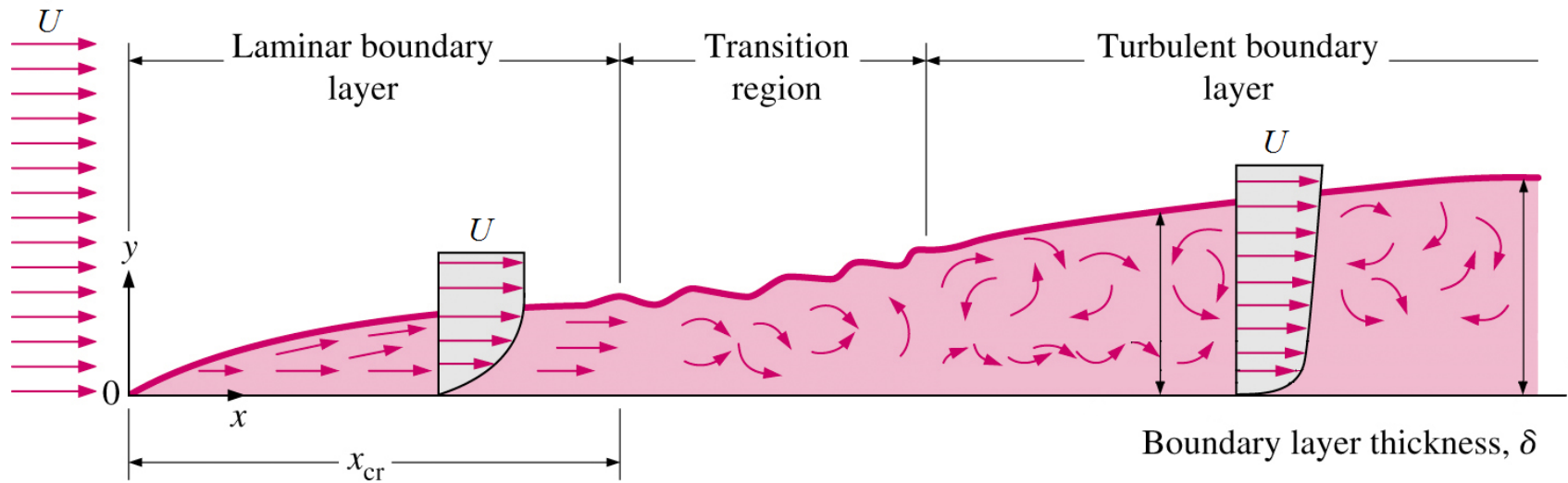
$$C_f = \frac{D_f}{\frac{1}{2}\rho U^2 A} \left(= \frac{1}{bL} \int_0^L \frac{2\tau_w(x)}{\rho U^2} b dx = \frac{1}{L} \int_0^L c_f(x) dx \right)$$

$$\therefore D_f = C_f \cdot \frac{1}{2}\rho U^2 A$$

Reynolds Number Regime

- Transition Reynolds number

$$Re_{x,tr} = 5 \times 10^5$$



Laminar boundary layer

- Prandtl/Blasius solution

$$\tau_w = 0.332U^{3/2} \sqrt{\frac{\rho\mu}{x}}$$

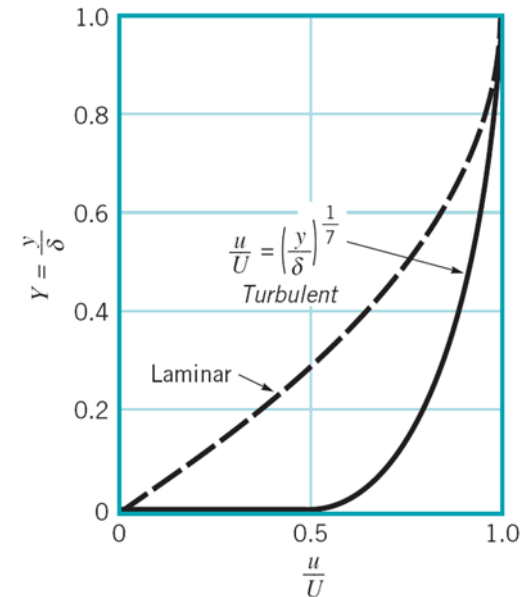
$$\frac{\delta(x)}{x} = \frac{5}{\sqrt{Re_x}} ; \quad c_f(x) = \frac{0.664}{\sqrt{Re_x}} ; \quad C_f = \frac{1.328}{\sqrt{Re_L}}$$

Turbulent boundary layer

- $\frac{u}{U} \approx \left(\frac{y}{\delta}\right)^{\frac{1}{7}}$ one – seventh – power law

- $c_f \approx 0.02 Re_{\delta}^{-\frac{1}{6}}$ power – law fit

- $\frac{\delta(x)}{x} = \frac{0.16}{Re_x^{\frac{1}{7}}}$; $c_f(x) = \frac{0.027}{Re_x^{\frac{1}{7}}}$; $C_f = \frac{0.031}{Re_L^{\frac{1}{7}}}$



- Valid for a fully turbulent flow over a smooth flat plate from the leading edge.
- Better results for sufficiently large $Re_L > 10^7$

Turbulent boundary layer – Contd.

- Alternate forms by using an experimentally determined shear stress formula:
- $$\tau_w = 0.0225\rho U^2 \left(\frac{\nu}{U\delta}\right)^{\frac{1}{4}}$$
- $$\frac{\delta(x)}{x} = 0.37Re_x^{-\frac{1}{5}}; \quad c_f(x) = \frac{0.058}{Re_x^{\frac{1}{5}}}; \quad C_f = \frac{0.074}{Re_L^{\frac{1}{5}}}$$
- Valid only in the range of the experimental data; $Re_L = 5 \times 10^5 \sim 10^7$ for smooth flat plate

Turbulent boundary layer – Contd.

- Other formulas for smooth flat plates are by using the logarithmic velocity-profile instead of the 1/7-power law:

$$\frac{\delta}{L} = c_f (0.98 \log Re_L - 0.732)$$

$$c_f = (2 \log Re_x - 0.65)^{-2.3}$$

$$C_f = \frac{0.455}{(\log_{10} Re_L)^{2.58}}$$

These formulas are valid in the whole range of $Re_L \leq 10^9$

Turbulent boundary layer – Contd.

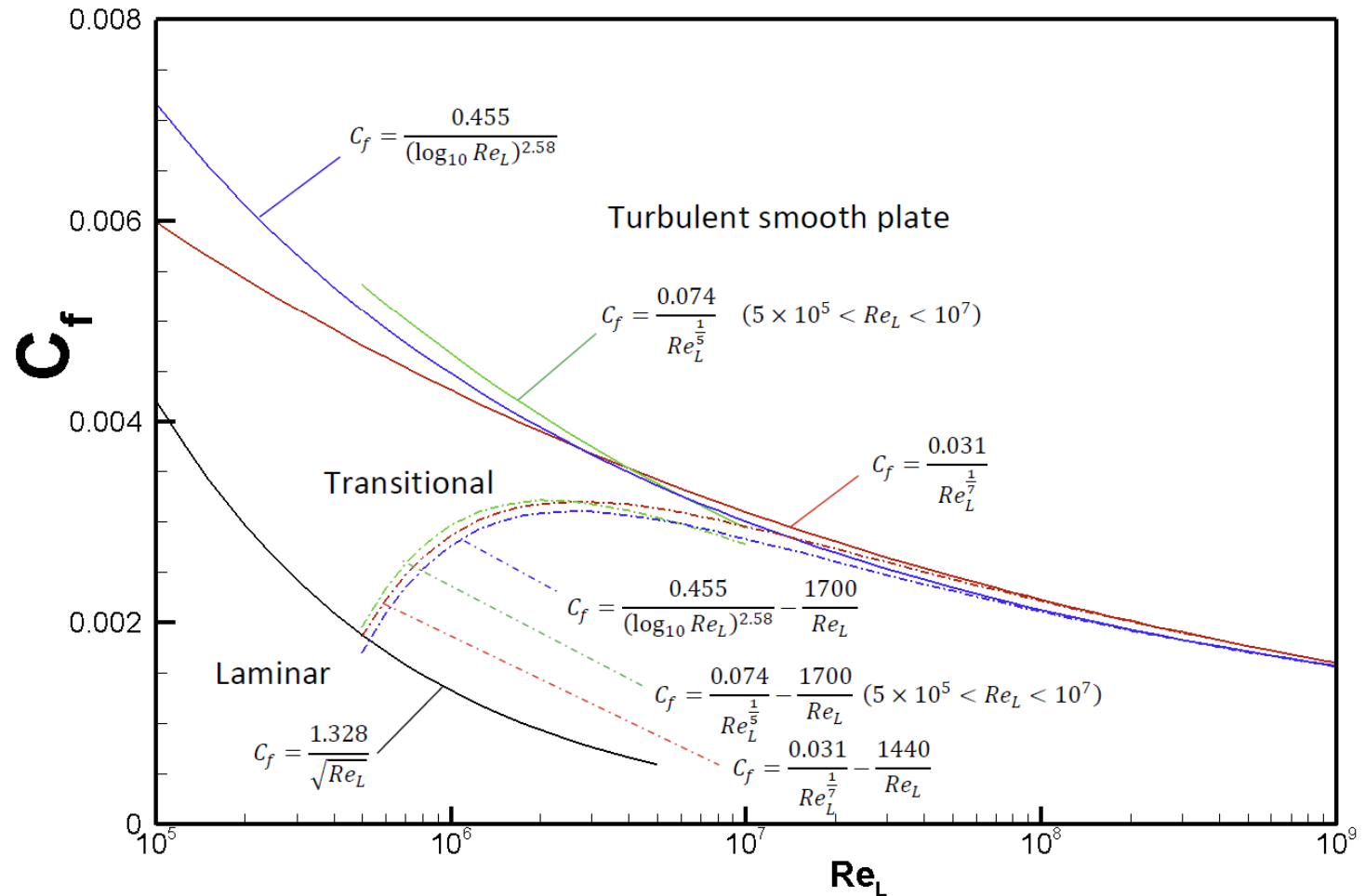
- Composite formulas (for flows initially laminar and subsequently turbulent with $Re_t = 5 \times 10^5$):

$$C_f = \frac{0.031}{Re_L^{\frac{1}{7}}} - \frac{1440}{Re_L}$$

$$C_f = \frac{0.074}{Re_L^{\frac{1}{5}}} - \frac{1700}{Re_L}$$

$$C_f = \frac{0.455}{(\log_{10} Re_L)^{2.58}} - \frac{1700}{Re_L}$$

Turbulent boundary layer – Contd.



Bluff Body Drag

- In general,

$$D = f(V, L, \rho, \mu, c, t, \varepsilon, \dots)$$

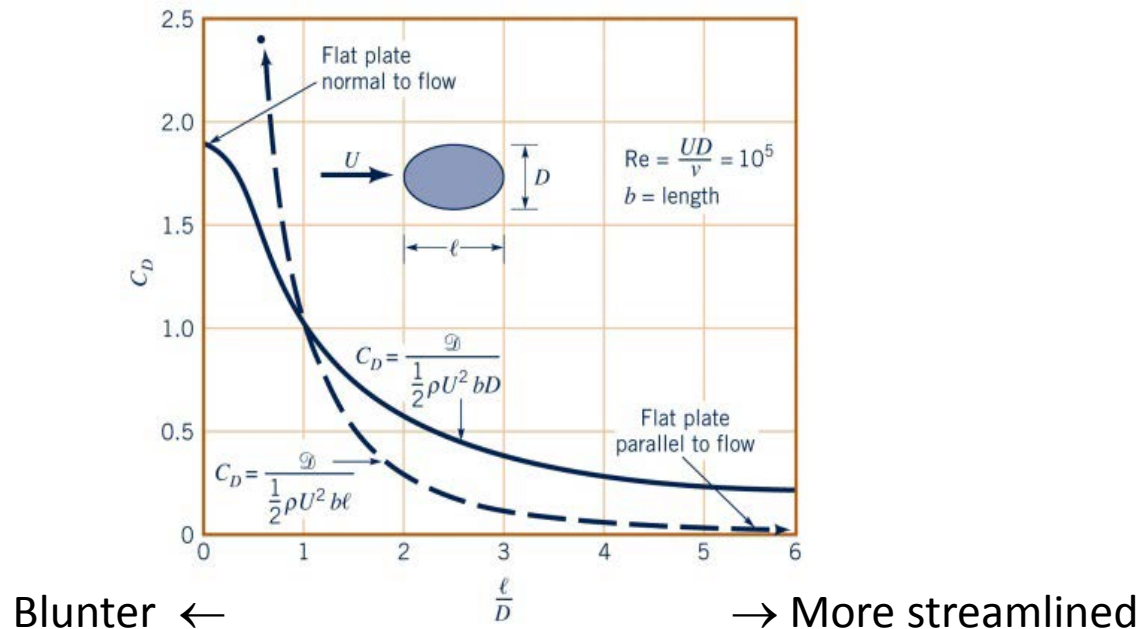
- Drag coefficient:

$$C_D = \frac{D}{\frac{1}{2}\rho V^2 A} = \phi \left(AR, \frac{t}{L}, Re, \frac{c}{V}, \frac{\varepsilon}{L}, \dots \right)$$

- For bluff bodies experimental data are used to determine C_D

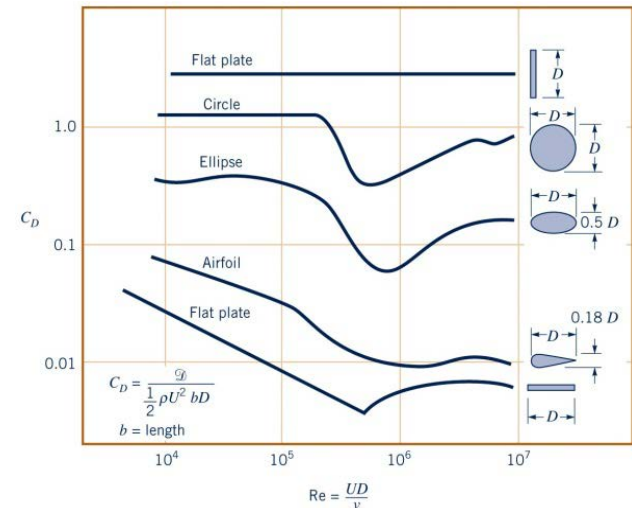
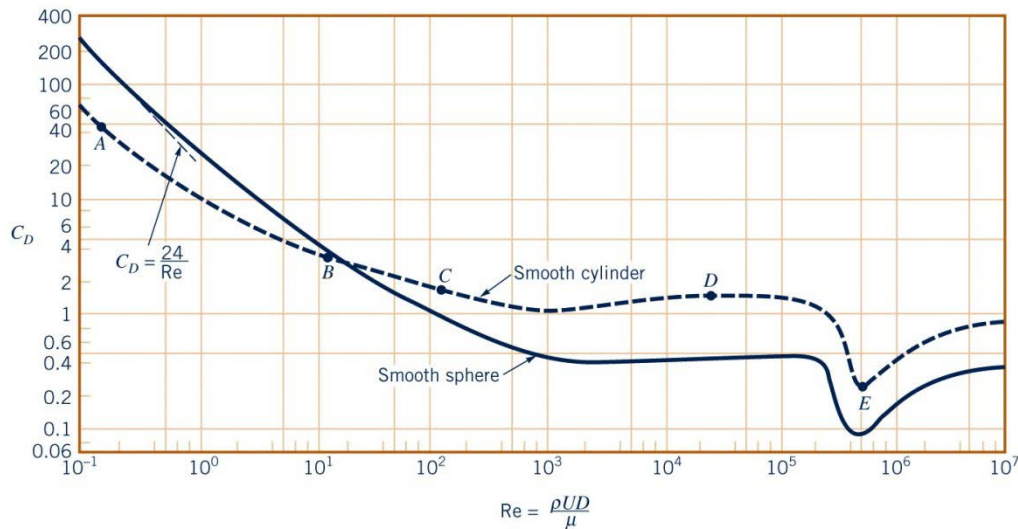
Shape dependence

- The blunter the body, the larger the drag coefficient
- The amount of streamlining can have a considerable effect



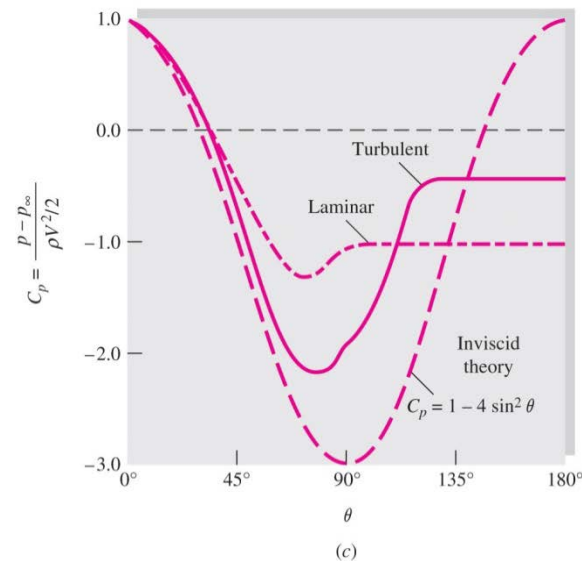
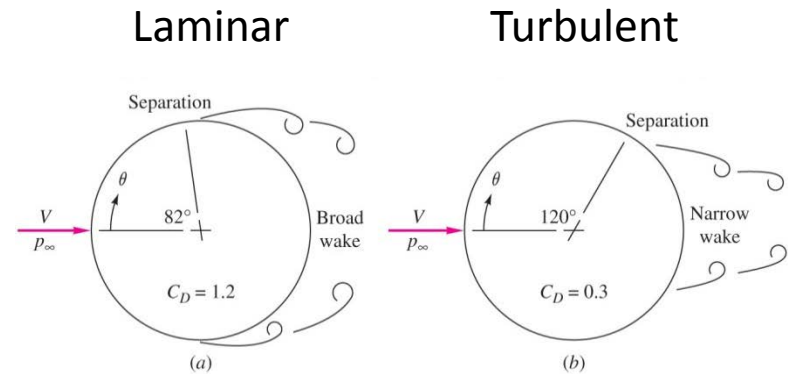
Reynolds number dependence

- Very low Re flow ($Re < 1$)
 - Inertia effects are negligible (creeping flow)
 - $C_D \sim Re^{-1}$
 - Streamlining can actually increase the drag (an increase in the area and shear force)
- Moderate Re flow ($10^3 < Re < 10^5$)
 - For streamlined bodies, $C_D \sim Re^{-\frac{1}{2}}$
 - For blunt bodies, $C_D \sim \text{constant}$
- Very large Re flow (turbulent boundary layer)
 - For streamlined bodies, C_D increases
 - For relatively blunt bodies, C_D decreases when the flow becomes turbulent ($10^5 < Re < 10^6$)
- For extremely blunt bodies, $C_D \sim \text{constant}$



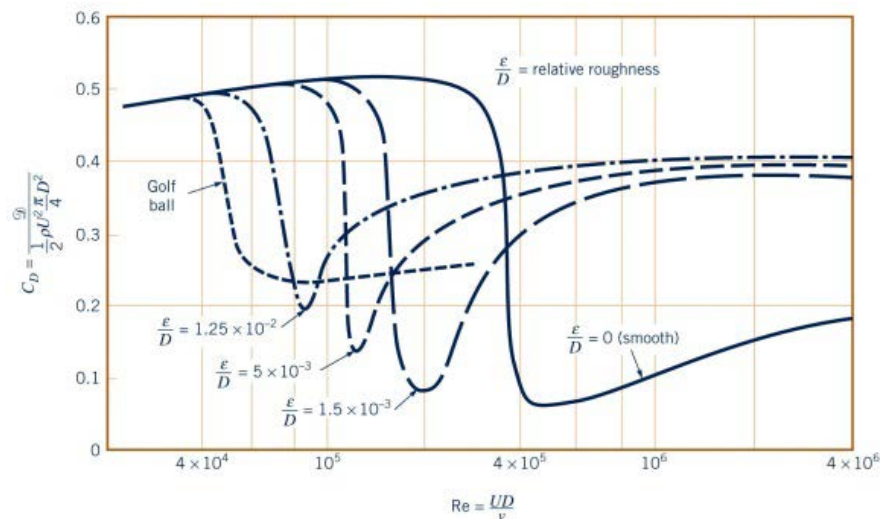
Separation

- Fluid stream detaches from a surface of a body at sufficiently high velocities.
- Only appears in viscous flows.
- Inside a separation region: low-pressure, existence of recirculating /backflows; viscous and rotational effects are the most significant



Surface roughness

- For streamlined bodies, the drag increases with increasing surface roughness
- For blunt bodies, an increase in surface roughness can actually cause a decrease in the drag.
- For extremely blunt bodies, the drag is independent of the surface roughness

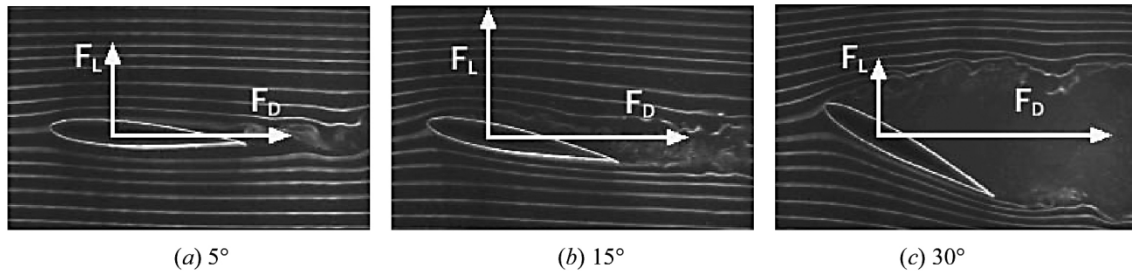


Lift

- Lift, L : Resultant force normal to the upstream velocity

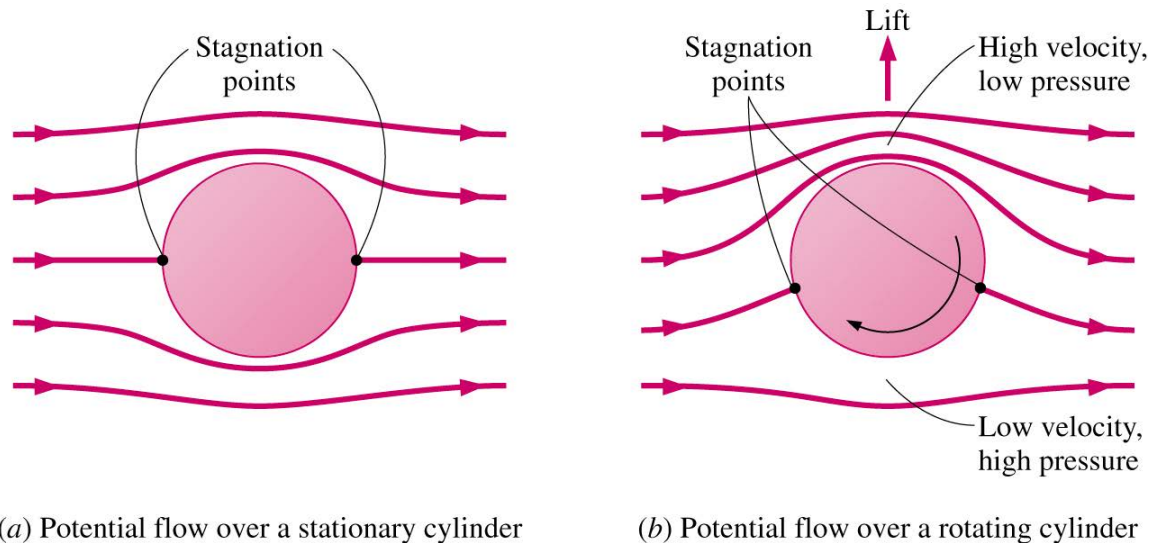
$$C_L = \frac{L}{\frac{1}{2}\rho U^2 A}$$

$$L = C_L \cdot \frac{1}{2}\rho U^2 A$$



Magnus Effect

- Lift generation by spinning
- Breaking the symmetry causes a lift



Minimum Flight Velocity

- Total weight of an aircraft should be equal to the lift

$$W = F_L = \frac{1}{2}C_{L,max}\rho V_{min}^2 A$$

Thus,

$$V_{min} = \sqrt{\frac{2W}{\rho C_{L,max} A}}$$