

Review for Exam2

11. 13. 2015

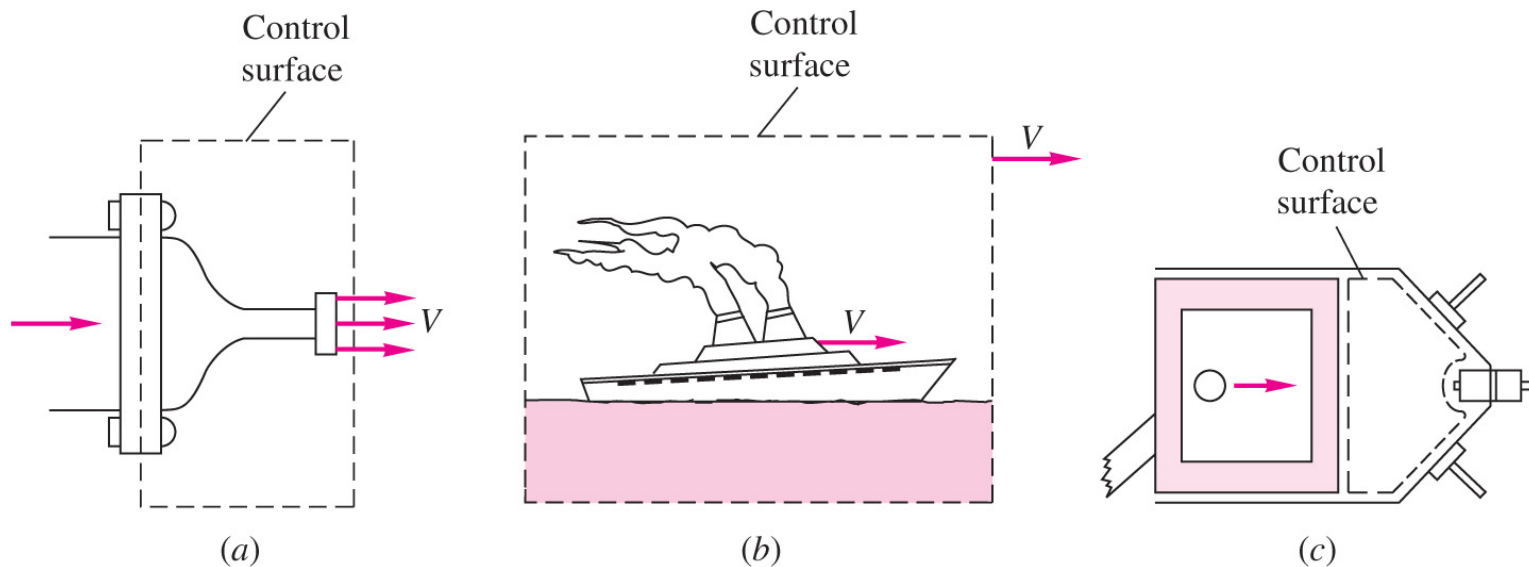
Hyunse Yoon, Ph.D.

Adjunct Assistant Professor
Department of Mechanical Engineering, University of Iowa

Assistant Research Scientist
IIHR-Hydroscience & Engineering, University of Iowa

System vs. Control volume

- **System:** A collection of real matter of fixed identity.
- **Control volume (CV):** A geometric or an imaginary volume in space through which fluid may flow. A CV may move or deform.



Laws of Mechanics for a System

Laws of mechanics are written for a system, i.e., for a fixed amount of matter

- Conservation of mass

$$\frac{Dm}{Dt} = 0$$

- Conservation of momentum

$$\frac{D(m\underline{V})}{Dt} = m\underline{a} = \underline{F}$$

Note:

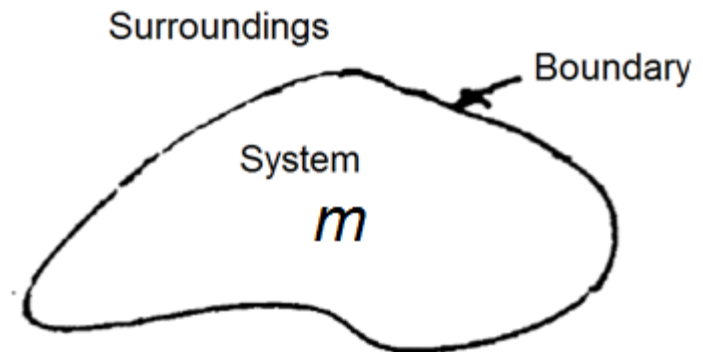
$$\frac{D(m\underline{V})}{Dt} = \underbrace{\frac{Dm}{Dt}}_{=0} \underline{V} + m \underbrace{\frac{D\underline{V}}{Dt}}_{=\underline{a}} = m\underline{a}$$

- Conservation of energy

$$\frac{DE}{Dt} = \dot{Q} - \dot{W}$$

Governing Differential Eq. (GDE):

$$\therefore \frac{D}{Dt} \underbrace{(m, m\underline{V}, E)}_{\substack{\text{system extensive} \\ \text{properties, } B_{\text{sys}}}} = \text{RHS}$$



Reynolds Transport Theorem (RTT)

- In fluid mechanics, we are usually interested in a region of space, i.e., CV and not particular systems. Therefore, we need to transform GDE's from a system to a CV, which is accomplished through the use of RTT

$$\underbrace{\frac{DB_{\text{sys}}}{Dt}}_{\text{time rate of change of } B \text{ for a system}} = \underbrace{\frac{D}{Dt} \int_{\text{CV}(\underline{x}, t)} \beta \rho dV}_{\text{time rate of change of } B \text{ in CV}} + \underbrace{\int_{\text{CS}(\underline{x}, t)} \beta \rho \underline{V}_R \cdot d\underline{A}}_{\text{net flux of } B \text{ across CS}}$$

where, $\beta = \frac{dB}{dm} = (1, \underline{V}, e)$ for $B = (m, m\underline{V}, E)$

- Fixed CV,

$$\frac{DB_{\text{sys}}}{Dt} = \frac{\partial}{\partial t} \int_{\text{CV}} \beta \rho dV + \int_{\text{CS}} \beta \rho \underline{V} \cdot d\underline{A}$$

Note:

$$B_{\text{CV}} = \int_{\text{CV}} \beta dm = \int_{\text{CV}} \beta \rho dV$$

$$\dot{B}_{\text{CS}} = \int_{\text{CS}} \beta d\dot{m} = \int_{\text{CS}} \beta \rho \underline{V} \cdot d\underline{A}$$

Continuity Equation

- RTT with $B = m$ and $\beta = 1$,

$$\frac{\partial}{\partial t} \int_{CV} \rho dV + \int_{CS} \rho \underline{V} \cdot d\underline{A} = 0$$

- Steady flow,

$$\int_{CS} \rho \underline{V} \cdot d\underline{A} = 0$$

- Simplified form,

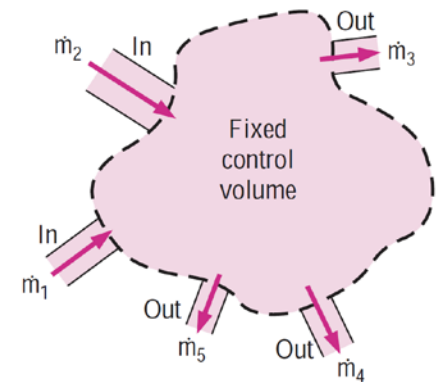
$$\sum \dot{m}_{out} - \sum \dot{m}_{in} = 0$$

- Conduit flow with one inlet (1) and one outlet (2):

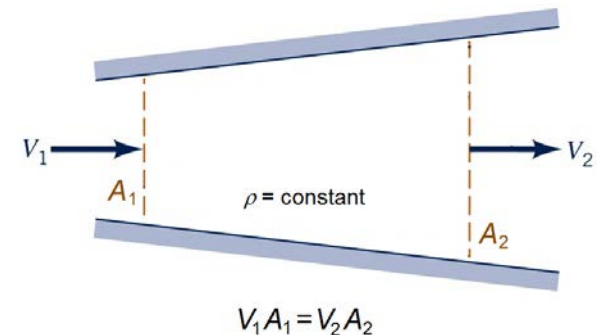
$$\rho_2 V_2 A_2 - \rho_1 V_1 A_1 = 0$$

If $\rho = \text{constant}$,

$$V_1 A_1 = V_2 A_2$$



Note: $\dot{m} = \rho Q = \rho V A$



Momentum Equation

- RTT with $B = m\underline{V}$ and $\beta = \underline{V}$,

$$\frac{\partial}{\partial t} \int_{CV} \underline{V} \rho dV + \int_{CS} \underline{V} \rho \underline{V} \cdot d\underline{A} = \underline{\Sigma F}$$

- Simplified form:

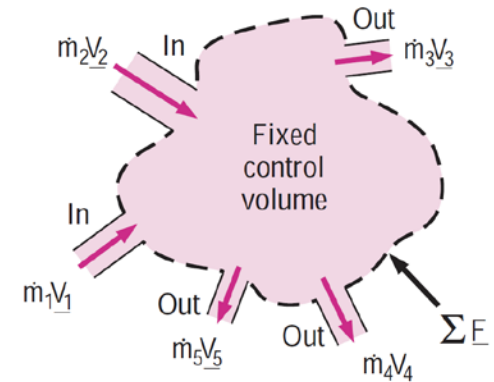
$$\underline{\Sigma(\dot{m}\underline{V})}_{out} - \underline{\Sigma(\dot{m}\underline{V})}_{in} = \underline{\Sigma F}$$

or in component forms,

$$\underline{\Sigma(\dot{m}u)}_{out} - \underline{\Sigma(\dot{m}u)}_{in} = \underline{\Sigma F}_x$$

$$\underline{\Sigma(\dot{m}v)}_{out} - \underline{\Sigma(\dot{m}v)}_{in} = \underline{\Sigma F}_y$$

$$\underline{\Sigma(\dot{m}w)}_{out} - \underline{\Sigma(\dot{m}w)}_{in} = \underline{\Sigma F}_z$$



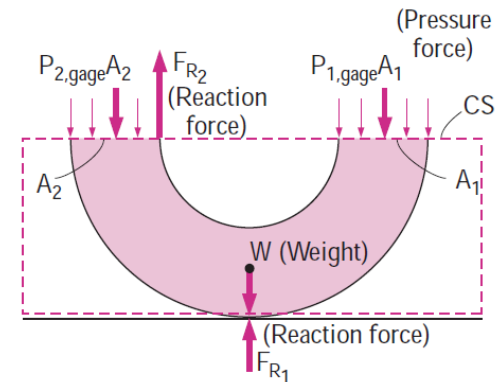
Note: If $\underline{V} = u\hat{i} + v\hat{j} + w\hat{k}$ is normal to CS, $\dot{m} = \rho VA$, where $V = |\underline{V}|$.

Momentum Equation – Contd.

- External forces:

$$\sum \underline{F} = \sum \underline{F}_{\text{body}} + \sum \underline{F}_{\text{surface}} + \sum \underline{F}_{\text{other}}$$

- $\sum \underline{F}_{\text{body}} = \sum \underline{F}_{\text{gravity}}$
 - $\sum \underline{F}_{\text{gravity}}$: gravity force (i.e., weight)
- $\sum \underline{F}_{\text{Surface}} = \sum \underline{F}_{\text{pressure}} + \sum \underline{F}_{\text{friction}} + \sum \underline{F}_{\text{other}}$
 - $\sum \underline{F}_{\text{pressure}}$: pressure forces normal to CS
 - $\sum \underline{F}_{\text{friction}}$: viscous friction forces tangent to CS
- $\sum \underline{F}_{\text{other}}$: anchoring forces or reaction forces

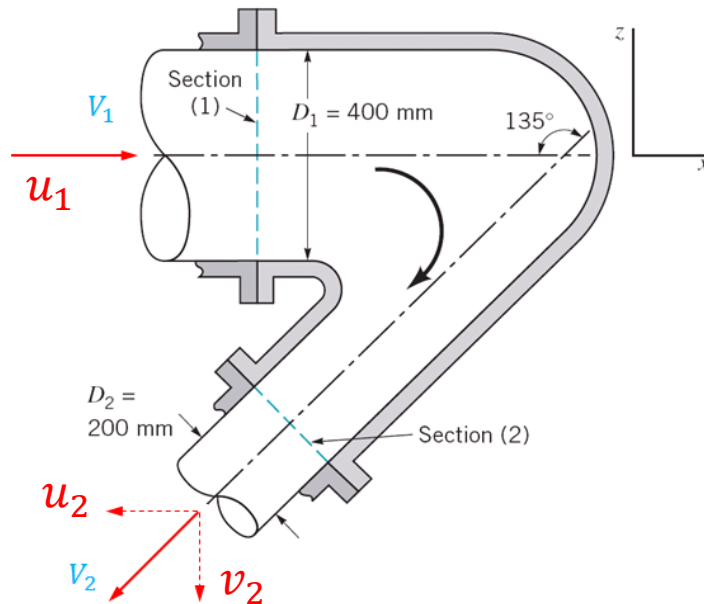


An 180° elbow supported by the ground

In most flow systems, the force \vec{F} consists of weights, pressure forces, and reaction forces. Gage pressures are used here since atmospheric pressure cancels out on all sides of the control surface.

Note: Surface forces arise as the CV is isolated from its surroundings, similarly to drawing a free-body diagram. A well-chosen CV exposes only the forces that are to be determined and a minimum number of other forces

Example (Bend)



Inlet (1):

$$\begin{aligned}\dot{m}_1 &= \rho V_1 A_1 \\ u_1 &= V_1 \\ v_1 &= 0\end{aligned}$$

Outlet (2):

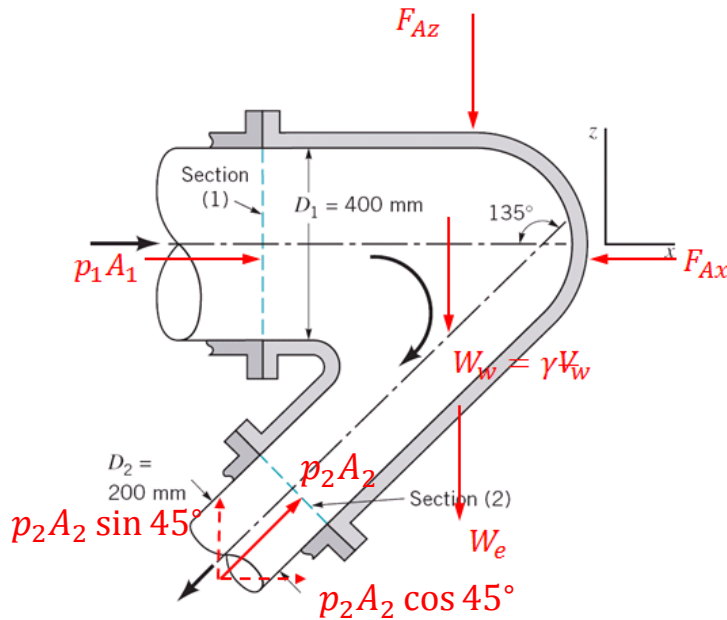
$$\begin{aligned}\dot{m}_2 &= \rho V_2 A_2 \\ u_2 &= -V_2 \cos 45^\circ \\ v_2 &= -V_2 \sin 45^\circ\end{aligned}$$

$$\begin{aligned}(\dot{m}u)_{\text{out}} - (\dot{m}u)_{\text{in}} &= (\rho V_2 A_2)(-V_2 \cos 45^\circ) - (\rho V_1 A_1)(V_1) \\ (\dot{m}v)_{\text{out}} - (\dot{m}v)_{\text{in}} &= (\rho V_2 A_2)(-V_2 \sin 45^\circ) - (\rho V_1 A_1)(0)\end{aligned}$$

Since $\rho V_1 A_1 = \rho V_2 A_2$,

$$\begin{aligned}(\dot{m}u)_{\text{out}} - (\dot{m}u)_{\text{in}} &= -(\rho V_2 A_2)(V_2 \cos 45^\circ + V_1) \\ (\dot{m}v)_{\text{out}} - (\dot{m}v)_{\text{in}} &= -\rho V_2^2 A_2 \sin 45^\circ\end{aligned}$$

Example – Contd.



$\Sigma F_x:$

- 1) Body force = 0
- 2) Pressure force = $p_1 A_1 + p_2 A_2 \cos 45^\circ$
- 3) Anchoring force = $-F_{Ax}$

$\Sigma F_y:$

- 1) Body force = $-W_w - W_e$
- 2) Pressure force = $p_2 A_2 \sin 45^\circ$
- 3) Anchoring force = $-F_{Az}$

Thus,

$$-(\rho V_2 A_2)(V_2 \cos 45^\circ + V_1) = p_1 A_1 + p_2 A_2 \cos 45^\circ - F_{Ax}$$

$$-\rho V_2^2 A_2 \sin 45^\circ = -\gamma V_w - W_e + p_2 A_2 \sin 45^\circ - F_{Az}$$

$$\therefore F_{Ax} = (\rho V_2 A_2)(V_2 \cos 45^\circ + V_1) + p_1 A_1 + p_2 A_2 \cos 45^\circ$$

$$F_{Az} = \rho V_2^2 A_2 \sin 45^\circ - \gamma V_w - W_e + p_2 A_2 \sin 45^\circ$$

Energy Equation

- RTT with $B = E$ and $\beta = e$,

$$\frac{\partial}{\partial t} \int_{CV} e \rho dV + \int_{CS} e \rho \underline{V} \cdot d\underline{A} = \dot{Q} - \dot{W}$$

- Simplified form:

$$\frac{p_{in}}{\gamma} + \alpha_{in} \frac{V_{in}^2}{2g} + z_{in} + h_p = \frac{p_{out}}{\gamma} + \alpha_{out} \frac{V_{out}^2}{2g} + z_{out} + h_t + h_L$$

- V in energy equation refers to average velocity \bar{V}
- α : kinetic energy correction factor = $\begin{cases} 1 & \text{for uniform flow across CS} \\ 2 & \text{for laminar pipe flow} \\ \approx 1 & \text{for turbulent pipe flow} \end{cases}$

Energy Equation - Contd.

Uniform flow across CS's:

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + z_1 + h_p = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + z_1 + h_t + h_L$$

- Pump head $h_p = \frac{\dot{W}_p}{\dot{m}g} = \frac{\dot{W}_p}{\rho Qg} = \frac{\dot{W}_p}{\gamma Q} \Rightarrow \dot{W}_p = \dot{m}gh_p = \rho g Q h_p = \gamma Q h_p$
- Turbine head $h_t = \frac{\dot{W}_t}{\dot{m}g} = \frac{\dot{W}_t}{\rho Qg} = \frac{\dot{W}_t}{\gamma Q} \Rightarrow \dot{W}_t = \dot{m}gh_t = \rho g Q h_t = \gamma Q h_t$
- Head loss $h_L = \text{loss}/g = (\hat{u}_2 - \hat{u}_1)/g - \dot{Q}/\dot{m}g > 0$

Example (Pump)

Energy equation:

$$\frac{p_1}{\gamma} + \frac{V_1^2}{2g} + z_1 + h_p = \frac{p_2}{\gamma} + \frac{V_2^2}{2g} + z_2 + h_t + h_L$$

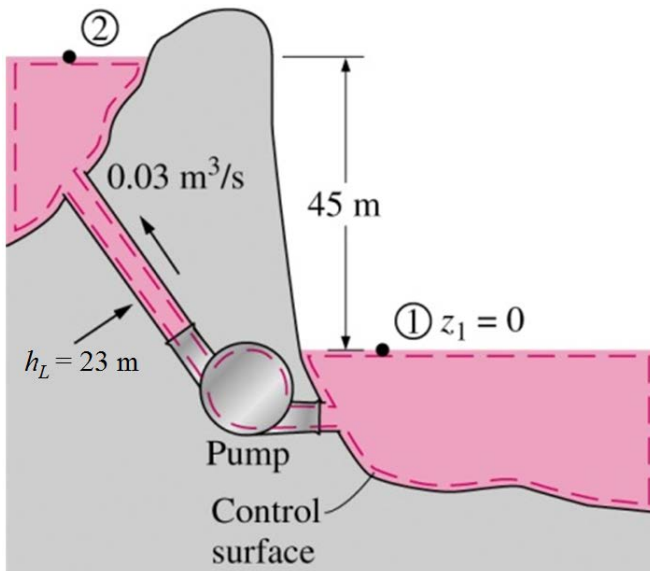
With $p_1 = p_2 = 0$, $V_1 = V_2 \approx 0$, $h_t = 0$, and $h_L = 23$ m

$$h_p = (z_2 - z_1) + h_L = 45 + 23 = 68 \text{ m}$$

Pump power,

$$\dot{W}_p = \gamma Q h_p = \frac{(68)(9790)(0.03)}{746} = 80 \text{ hp}$$

(Note: 1 hp = 746 N·m/s = 550 ft·lbf/s)



Differential Analysis

A microscopic description of fluid motions for a fluid particle by using differential equations*,

- Continuity equation

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \underline{V}) = 0$$

- Momentum equation

$$\rho \left(\frac{\partial \underline{V}}{\partial t} + \underline{V} \cdot \nabla \underline{V} \right) = \rho \underline{g} - \nabla p + \nabla \cdot \tau_{ij}$$

*CV analysis is a macroscopic description of fluid motions by using integral equations (RTT).

Navier-Stokes Equations

For incompressible, Newtonian fluids,

- Continuity:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$

- Momentum:

$$\rho \left(\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) = -\frac{\partial p}{\partial x} + \rho g_x + \mu \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right)$$

$$\rho \left(\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} \right) = -\frac{\partial p}{\partial y} + \rho g_y + \mu \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right)$$

$$\rho \left(\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} \right) = -\frac{\partial p}{\partial z} + \rho g_z + \mu \left(\frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} \right)$$

Exact Solutions of NS Eqns.

The flow of interest is assumed additionally (than incompressible & Newtonian), for example,

- 1) Steady (i.e., $\partial/\partial t = \mathbf{0}$ for any variable)
- 2) Parallel such that the y -component of velocity is zero (i.e., $v = \mathbf{0}$)
- 3) Purely two dimensional (i.e., $w = \mathbf{0}$ and $\partial/\partial z = \mathbf{0}$ for any velocity component)
- 4) Fully developed (i.e., $\partial/\partial x = \mathbf{0}$ for any velocity component)

e.g.)

$$\rho \left[\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + \tilde{v} \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right] = -\frac{\partial p}{\partial x} + \rho g_x + \mu \left[\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right]$$

or

$$\mu \frac{d^2 u}{dy^2} = \frac{\partial p}{\partial x} - \rho g_x$$

Boundary Conditions

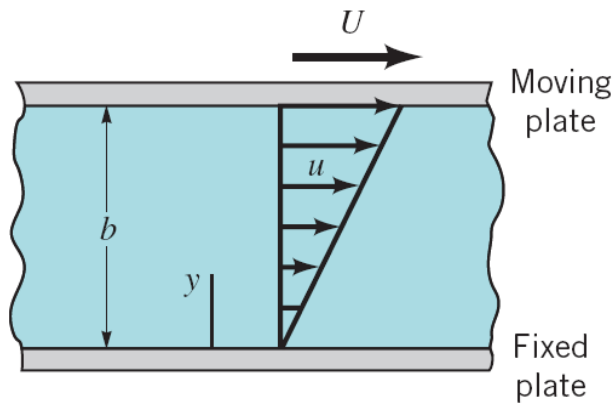
Common BC's:

- **No-slip condition** ($\underline{V}_{\text{fluid}} = \underline{V}_{\text{wall}}$; for a stationary wall $\underline{V}_{\text{fluid}} = 0$)
- Interface boundary condition ($\underline{V}_A = \underline{V}_B$ and $\tau_{s,A} = \tau_{s,B}$)
- Free-surface boundary condition ($p_{\text{liquid}} = p_{\text{gas}}$ and $\tau_{s,\text{liquid}} = 0$)

Other BC's:

- Inlet/outlet boundary condition
- Symmetry boundary condition
- Initial condition (for unsteady flow problem)

Example: No pressure gradient



$$\mu \frac{d^2 u}{dy^2} = 0$$

Integrate twice,

$$u(y) = C_1 y + C_2$$

B.C.,

$$u(0) = (C_1)(0) + C_2 = 0 \Rightarrow C_2 = 0$$

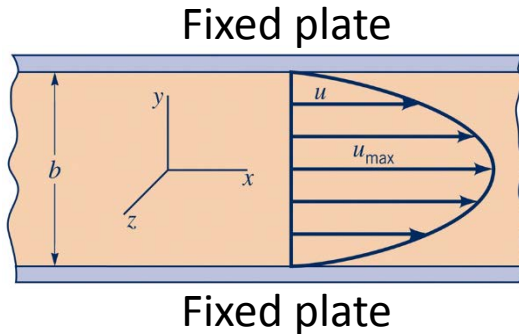
$$u(b) = (C_1)(b) + C_2 = U \Rightarrow C_1 = \frac{U}{b}$$

$$\therefore u(y) = \frac{U}{b} y$$

Analysis:

$$\tau_w = \mu \left. \frac{du}{dy} \right|_{y=0} = (\mu) \left(\frac{U}{b} \right) = \frac{\mu U}{b}$$

Example: with Pressure Gradient



$$\mu \frac{d^2 u}{dy^2} = \frac{dp}{dx}$$

Integrate twice,

$$u(y) = \frac{1}{2\mu} \frac{dp}{dx} y^2 + C_1 y + C_2$$

B.C.,

$$u(0) = \left(\frac{1}{2\mu} \frac{dp}{dx} \right) (0)^2 + (C_1)(0) + C_2 = 0 \Rightarrow C_2 = 0$$

$$u(b) = \left(\frac{1}{2\mu} \frac{dp}{dx} \right) (b)^2 + (C_1)(b) + C_2 = 0 \Rightarrow C_1 = -\frac{1}{2\mu} \frac{dp}{dx} b$$

$$\therefore u(y) = \frac{1}{2\mu} \left(\frac{dp}{dx} \right) (y^2 - by)$$

Analysis:

$$q = \int_{-h}^h u dy = -\frac{b^3}{12\mu} \left(\frac{\partial p}{\partial x} \right)$$

$$\tau_w = \mu \left. \frac{du}{dy} \right|_{y=0} = -\frac{b}{2} \left(\frac{\partial p}{\partial x} \right)$$

Example: Inclined wall

$$\mu \frac{d^2 u}{dy^2} = -\rho g_x$$

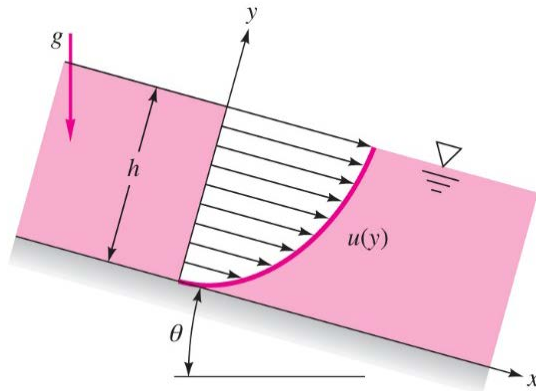
Integrate twice,

$$u(y) = -\frac{\rho g_x}{2\mu} y^2 + C_1 y + C_2$$

B.C.,

$$u(0) = \left(-\frac{\rho g_x}{\mu}\right)(0)^2 + (C_1)(0) + C_2 = 0 \Rightarrow C_2 = 0$$

$$\left(\frac{du}{dy}\right)_{y=h} = \left(-\frac{\rho g_x}{\mu}\right)(h) + C_1 = 0 \Rightarrow C_1 = \frac{\rho g_x}{\mu} h$$



Note:

$$\underline{g} = g_x \hat{i} + g_y \hat{j}$$

where,

$$g_x = g \sin \theta$$

$$g_y = -g \cos \theta$$

$$\therefore u(y) = \frac{\rho g_x}{\mu} \left(hy - \frac{y^2}{2} \right)$$

Analysis:

$$q = \int_0^h u dy = \frac{\rho g_x}{\mu} \frac{h^3}{3}$$

$$\tau_w = \mu \left(\frac{du}{dy}\right)_{y=0} = (\mu) \left(\frac{\rho g_x}{\mu} h\right) = \rho g_x h$$

Buckingham Pi Theorem

- For any physically meaningful equation involving **n variables**, such as

$$u_1 = f(u_2, u_3, \dots, u_n)$$

with minimum number of **m reference dimensions**, the equation can be rearranged into product of **r dimensionless pi terms**.

$$\Pi_1 = \phi(\Pi_2, \Pi_3, \dots, \Pi_r)$$

where,

$$r = n - m$$

Similarity and Model Testing

If all relevant dimensionless parameters have the same corresponding values for model and prototype, flow conditions for a model test are completely similar to those for prototype.

For,

$$\Pi_1 = \phi(\Pi_2, \dots, \Pi_n)$$

Similarity requirements:

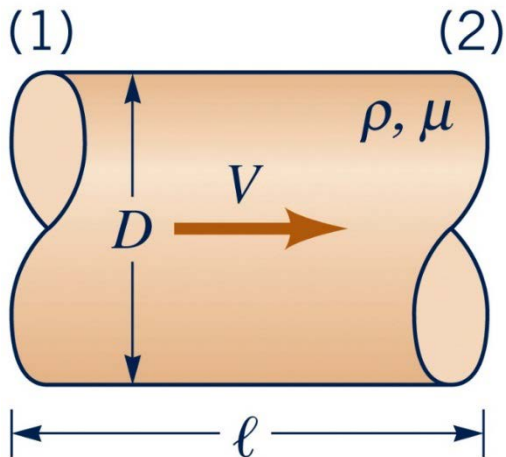
$$\begin{aligned}\Pi_{2,\text{model}} &= \Pi_{2,\text{prototype}} \\ &\vdots \\ \Pi_{n,\text{model}} &= \Pi_{n,\text{prototype}}\end{aligned}$$

Prediction equation:

$$\Pi_{1,\text{model}} = \Pi_{1,\text{prototype}}$$

Example (Repeating Variable Method)

Example: The pressure drop per unit length Δp_ℓ in a pipe flow is a function of the pipe diameter D and the fluid density ρ , viscosity μ , and velocity V .



$$\Delta p_\ell = (p_1 - p_2)/\ell$$

$$\Delta p_\ell = f(D, \rho, \mu, V)$$

Δp_ℓ	D	ρ	μ	V
$\{ML^{-2}T^{-2}\}$	$\{L\}$	$\{ML^{-3}\}$	$\{ML^{-1}T^{-1}\}$	$\{LT^{-1}\}$

$$r = n - m = 5 - 3 = 2$$

Example – Contd.

Select $m = 3$ repeating variables, (D, V, ρ) for (L, T, M) , then

$$\Pi_1 = D^a V^b \rho^c \Delta p_\ell \doteq (L)^a (LT^{-1})^b (ML^{-3})^c (ML^{-2}T^{-2}) \doteq M^0 L^0 T^0$$

$$\begin{aligned} a + b - 3c - 2 &= 0 \\ -b - 2 &= 0 \\ c + 1 &= 0 \end{aligned} \Rightarrow a = -1, b = -2, c = -1$$

$$\Rightarrow \Pi_1 = D^{-1} V^{-2} \rho^{-1} \Delta p_\ell = \frac{\Delta p_\ell D}{\rho V^2}$$

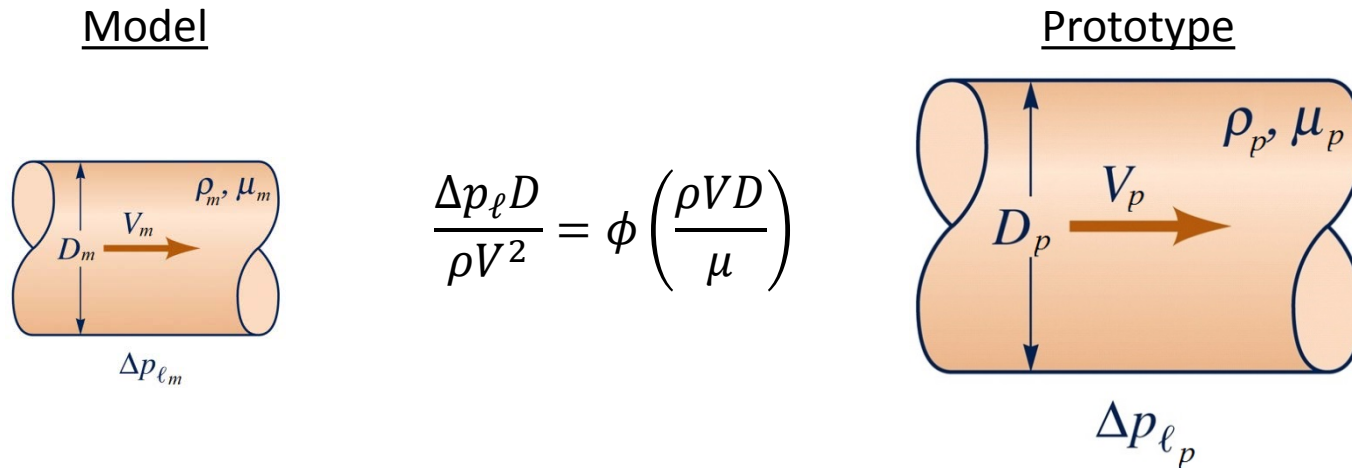
$$\Pi_2 = D^a V^b \rho^c \mu \doteq (L)^a (LT^{-1})^b (ML^{-3})^c (ML^{-1}T^{-1}) \doteq M^0 L^0 T^0$$

$$\begin{aligned} a + b - 3c - 1 &= 0 \\ -b - 1 &= 0 \\ c + 1 &= 0 \end{aligned} \Rightarrow a = -1, b = -1, c = -1$$

$$\Rightarrow \Pi_2 = D^{-1} V^{-1} \rho^{-1} \mu = \frac{\mu}{DV\rho}$$

$$\therefore \frac{\Delta p_\ell D}{\rho V^2} = \phi \left(\frac{\rho V D}{\mu} \right)$$

Example (Model Testing)



If,

$$\frac{\rho_m V_m D_m}{\mu_m} = \frac{\rho_p V_p D_p}{\mu_p} \quad (\text{similarity requirement})$$

Then,

$$\frac{\Delta p_{\ell m} D_m}{\rho_m V_m^2} = \frac{\Delta p_{\ell p} D_p}{\rho_p V_p^2} \quad (\text{Prediction equation})$$

Example – Contd.

Model (in water)

- $D_m = 0.1 \text{ m}$
- $\rho_m = 998 \text{ kg/m}^3$
- $\mu_m = 1.12 \times 10^{-3} \text{ N}\cdot\text{s/m}^2$
- $V_m = ?$
- $\Delta p_{\ell m} = 27.6 \text{ Pa/m}$

Prototype (in air)

- $D_p = 1 \text{ m}$
- $\rho_p = 1.23 \text{ kg/m}^3$
- $\mu_p = 1.79 \times 10^{-5} \text{ N}\cdot\text{s/m}^2$
- $V_p = 10 \text{ m/s}$
- $\Delta p_{\ell m} = ?$

Similarity requirement:

$$V_m = \left(\frac{\rho_p}{\rho_m}\right) \left(\frac{\mu_m}{\mu_p}\right) \left(\frac{D_p}{D_m}\right) V_p = \left(\frac{1.23}{998}\right) \left(\frac{1.12 \times 10^{-3}}{1.79 \times 10^{-5}}\right) \left(\frac{1}{0.1}\right) (10) = \mathbf{7.71 \text{ m/s}}$$

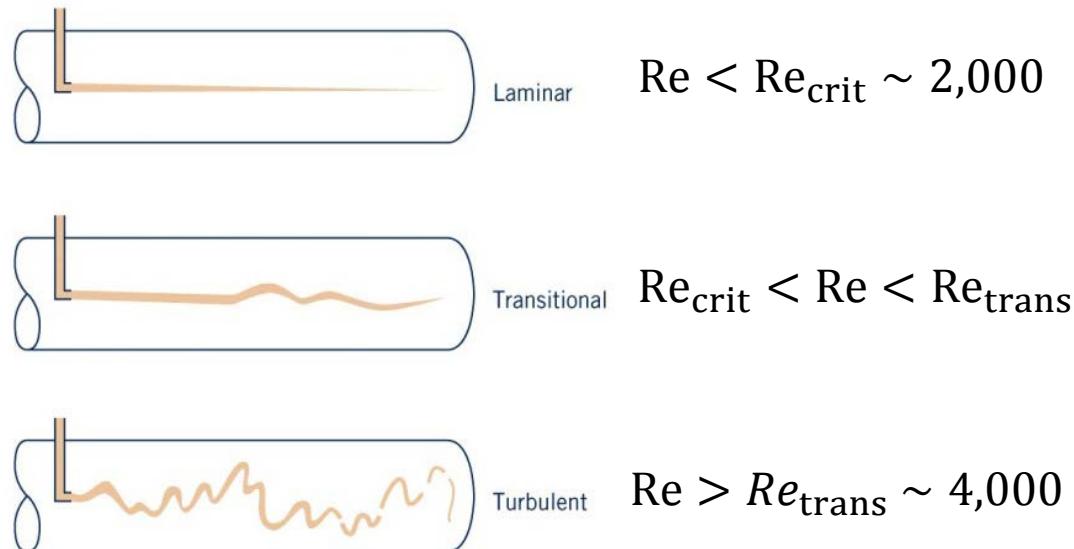
Prediction equation:

$$\Delta p_{\ell p} = \left(\frac{D_m}{D_p}\right) \left(\frac{\rho_p}{\rho_m}\right) \left(\frac{V_p}{V_m}\right)^2 \Delta p_{\ell m} = \left(\frac{0.1}{1}\right) \left(\frac{1.23}{998}\right) \left(\frac{10}{7.71}\right)^2 (27.6) = \mathbf{5.72 \times 10^{-3} \text{ Pa/m}}$$

Pipe Flow: Laminar vs. Turbulent

- Reynolds number regimes

$$Re = \frac{\rho V D}{\mu}$$



Flow in Pipes

- Basic piping problems:
 - Given the desired flow rate, what pressure drop (e.g., pump power) is needed to drive the flow (i.e., to overcome the head loss through piping)?
 - Given the pressure drop (e.g., pump power) available, what flow rate will ensue?
 - Given the pressure drop and the flow rate desired, what pipe diameter is needed?

Head Loss

$$h_L = h_{L \text{ major}} + h_{L \text{ minor}}$$

- $h_{L \text{ major}}$ (or h_f): Major loss, the loss due to viscous effects
- $h_{L \text{ minor}}$: Minor loss, the loss in the various pipe components

Darcy-Weisbach equation

$$h_f = f \frac{L V^2}{D 2g}$$

- $f = \frac{8\tau_w}{\rho V^2}$: Friction factor
- L : Pipe length
- D : Pipe diameter
- V : Average flow velocity across the pipe cross-section

Laminar Pipe Flow

- Exact solution exists by solving the NS equation

$$u(r) = V_{\max} \left[1 - \left(\frac{r}{R} \right)^2 \right], \quad V_{\max} = 2V$$

- Wall shear stress

$$\tau_w = -\mu \left. \frac{du}{dr} \right|_{r=R} = \frac{8\mu V}{D}$$

- Friction factor

$$f = \frac{8\tau_w}{\rho V^2} = \frac{64\mu}{\rho D V} = \frac{64}{\text{Re}}$$

- Head loss

$$h_f = f \frac{L V^2}{D 2g} = \frac{32\mu L V}{\gamma D^2} \quad \left(= \frac{128\mu L Q}{\pi \gamma D^4} \right)$$

Notes:

$$Q = VA$$

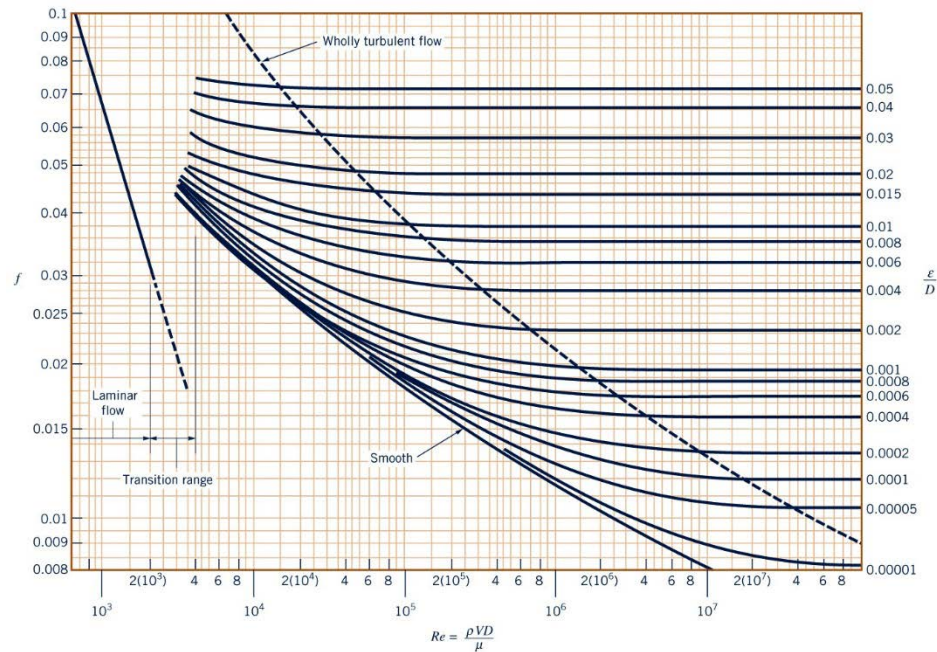
$$A = \frac{\pi D^2}{4}$$

Turbulent Pipe Flow

- From a dimensional analysis

$$f = \phi(\text{Re}, \varepsilon/D)$$

- Moody chart: Empirical functional dependency of f on Re and ε/D



Turbulent Pipe Flow – Cond.

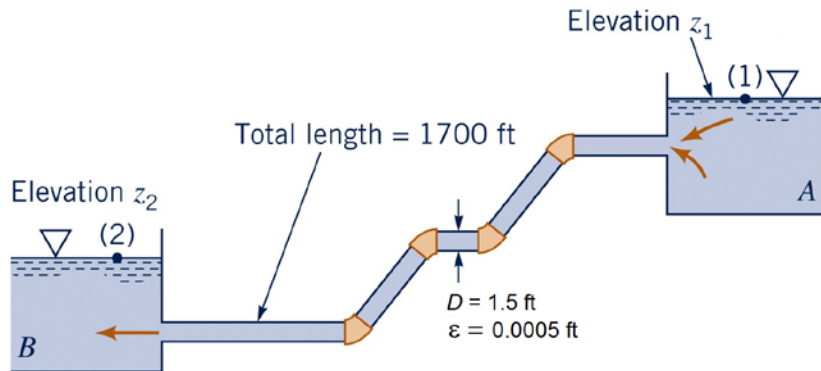
- Colebrook equation (difficult in its use as implicit)

$$\frac{1}{\sqrt{f}} = -2 \log \left(\frac{\varepsilon/D}{3.7} + \frac{2.51}{Re\sqrt{f}} \right)$$

- Haaland equation (easier to use as explicit but approximation)

$$\frac{1}{\sqrt{f}} = -1.8 \log \left[\left(\frac{\varepsilon/D}{3.7} \right)^{1.1} + \frac{6.9}{Re} \right]$$

Example (pipe flow)



If $D = 1.5$ ft and $Q = 25$ ft³/s, $\Delta z = z_1 - z_2$?
Neglect minor losses.

$$V = \frac{Q}{A} = \frac{25}{(\pi)(1.5)^2/4} = 14.1 \text{ ft/s}$$

$$\text{Re} = \frac{VD}{\nu} = \frac{(14.1)(1.5)}{1.21 \times 10^{-5}} = 1.75 \times 10^6 \text{ (turbulent)}$$

$$\epsilon/D = 0.0005/1.63 = 0.00033$$

Energy equation:

$$\frac{p_1}{\gamma} + \alpha_1 \frac{V_1^2}{2g} + z_1 = \frac{p_2}{\gamma} + \alpha_2 \frac{V_2^2}{2g} + z_2 + h_f$$

Since $p_1 = p_2 = 0$ and $V_1 = V_2$,

$$h_f = z_1 - z_2 = \Delta z$$

Friction factor,

$$\frac{1}{\sqrt{f}} = -1.8 \log \left[\left(\frac{0.00033}{3.7} \right)^{1.1} + \frac{6.9}{1.75 \times 10^6} \right] \Rightarrow f = 0.0159$$

Head loss

$$h_f = f \frac{L V^2}{D 2g} = (0.0159) \frac{(1700) (14.1)^2}{(1.5) (2)(32.2)} = 56 \text{ ft}$$

$$\therefore \Delta z = \mathbf{56 \text{ ft}}$$