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Fluid Mechanics

Class Notes Fall 2013

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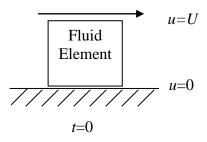
CHAPTER 1: INTRODUCTION AND BASIC CONCEPTS

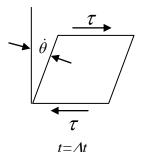
Fluids and the no-slip condition

Fluid mechanics is the science and technology of fluids either at rest (fluid statics) or in motion (fluid dynamics) and their effects on boundaries such as solid surfaces or interfaces with other fluids.

<u>Definition of a fluid:</u> A substance that deforms continuously when subjected to a shear stress

Consider a fluid between two parallel plates, which is subjected to a shear stress due to the impulsive motion of the upper plate





No slip condition: no relative motion between fluid and boundary, i.e., fluid in contact with lower plate is stationary, whereas fluid in contact with upper plate moves at speed *U*.

Fluid deforms, i.e., undergoes rate of strain θ due to shear stress τ

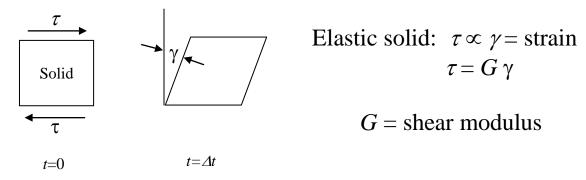
Newtonian fluid:

$$\tau \propto \dot{\theta} = \text{rate of strain}$$

$$\tau = \mu \dot{\theta}$$

 μ = coefficient of viscosity

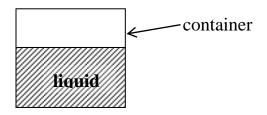
Such behavior is different from solids, which resist shear by static deformation (up to elastic limit of material)



Both liquids and gases behave as fluids

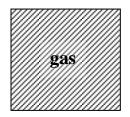
Liquids:

Closely spaced molecules with large intermolecular forces Retain volume and take shape of container



Gases:

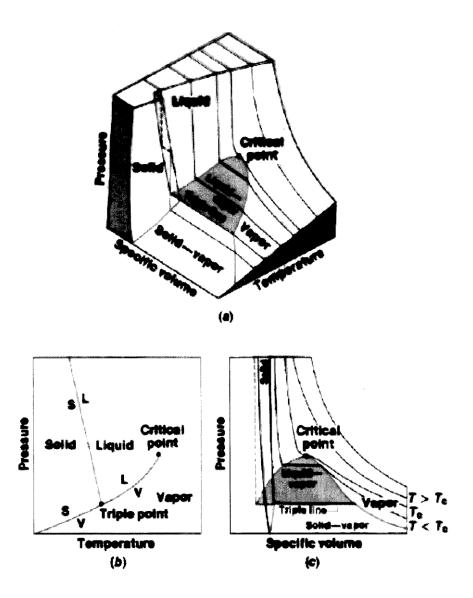
Widely spaced molecules with small intermolecular forces Take volume and shape of container



Recall *p-v-T* diagram from thermodynamics:

single phase, two phase, triple point (point at which solid, liquid, and vapor are all in equilibrium), critical point (maximum pressure at which liquid and vapor are both in equilibrium).

Liquids, gases, and two-phase liquid-vapor behave as fluids.



Continuum Hypothesis

In this course, the assumption is made that the fluid behaves as a continuum, i.e., the number of molecules within the smallest region of interest (a point) are sufficient that all fluid properties are point functions (single valued at a point).

For example:

Consider definition of density ρ of a fluid

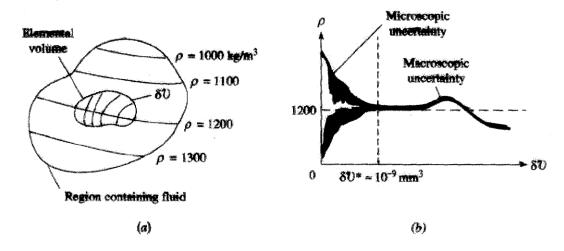
$$\rho(\underline{x},t) = \frac{\lim}{\delta \mathbf{V} \to \delta \mathbf{V}^*} \frac{\delta m}{\delta \mathbf{V}} \qquad \underline{\mathbf{x}} = \text{position vector} = x\mathbf{i} + y\mathbf{j} + z\mathbf{k}$$
$$\mathbf{t} = \text{time}$$

 δV^* = limiting volume below which molecular variations may be important and above which macroscopic variations may be important

 $\delta V^* \approx 10^{-9} \text{ mm}^3$ for all liquids and for gases at atmospheric pressure

 10^{-9} mm³ air (at standard conditions, 20° C and 1 atm) contains $3x10^{7}$ molecules such that $\delta M/\delta V = constant = \rho$

Note that typical "smallest" measurement volumes are about $10^{-3} - 10^0 \, \text{mm}^3 >> \delta \text{V}^*$ and that the "scale" of macroscopic variations are very problem dependent



²One atmosphere equals $2116 \text{ lbf/ft}^2 = 101,300 \text{ Pa}$.

Exception: rarefied gas flow

 $\delta \forall^*$ defines a point in the fluid, i.e., a fluid particle or infinitesimal material element used for deriving governing differential equations of fluid dynamics and at which all fluid properties are point functions:

 $\delta \forall^* >>$ molecular scales

 $\lambda = \text{mean free path} = 6 \times 10^{-8} \text{ m}$

 $t_{\lambda} = 10^{-10} \, s = time between collisions$

 $\delta \forall^* << \text{fluid motions scales} = 10^{-18} \text{ m}^3, \text{ i.e., } 1^* = 10^{-6} \text{ m}$ For laminar flow: $l_{max} \approx$ smallest geometry scales of the flow

 $U_{max} < U$ transition to turbulent flow

For turbulent flow: l_{max} and U_{max} determined by Kolmogorov scales at which viscous dissipation takes place, which for typical ship/airplane,

 $\eta \approx 2 \times 10^{-5}$ m (ship)/2.3×10⁻⁵ m (airplane)

 $u_n \approx 0.05 \text{ m/s (ship)}/1.64 \text{ m/s (airplane)}$

 $t_n \approx 4 \times 10^{-4} \text{ s (ship)}/1.4 \times 10^{-5} \text{ s (airplane)}$

Properties of Fluids

Fluids are characterized by their properties such as viscosity μ and density ρ , which we have already discussed with reference to definition of shear stress $\tau = \mu \dot{\theta}$ and the continuum hypothesis.

Properties can be both dimensional (i.e., expressed in either SI or BG units) or non-dimensional.

- Figure B.1 Dynamic (absolute) viscosity of common fluids as a function of temperature.
- Figure B.2 Kinematic viscosity of common fluids (at atmospheric pressure) as a function of temperature.
- Table B.1 Physical Properties of Water (BG Units)
- Table B.2 Physical Properties of Water (SI Units)
- Table B.3 Physical Properties of Air at Standard Atmospheric Pressure (BG Units)
- Table B.4 Physical Properties of Air at Standard Atmospheric Pressure (SI Units)
- Table 1.5 Approximate Physical Properties of Some Common Liquids (BG Units)
- Table 1.6 Approximate Physical Properties of Some Common Liquids (SI Units)
- Table 1.7 Approximate Physical Properties of Some Common Gases at Standard Atmospheric Pressure (BG Units)
- Table 1.8 Approximate Physical Properties of Some Common Gases at Standard Atmospheric Pressure (SI Units)

Basic Units

System International and British Gravitational Systems

| Primary Units | SI | BG |
|---------------|---------|--------------|
| Mass M | kg | slug=32.2lbm |
| Length L | m | ft |
| Time t | S | S |
| Temperature T | °C (°K) | °F (°R) |

Temperature Conversion:

$$^{\circ}$$
K = $^{\circ}$ C + 273

$$^{\circ}R = ^{\circ}F + 460$$

°K and °R are absolute scales, i.e., 0 at absolute zero. Freezing point of water is at 0°C and 32°F.

| Secondary | | | |
|-----------------------|------------------|------------------------|----------------------|
| (derived) units | Dimension | SI | BG |
| velocity <u>V</u> | L/t | m/s | ft/s |
| acceleration <u>a</u> | L/t^2 | m/s ² | ft/s ² |
| force <u>F</u> | ML/t^2 | $N (kg \cdot m/s^2)$ | lbf |
| pressure p | F/L ² | Pa (N/m ²) | lbf/ft ² |
| density ρ | M/L^3 | kg/m ³ | slug/ft ³ |
| internal energy u | FL/M | J/kg (N·m/kg) | BTU/lbm |

Table 1.3 Conversion Factors from BG and EE Units to SI Units.

Table 1.4 Conversion Factors from SI Units to BG and EE Units.

Weight and Mass

 $\underline{F} = m\underline{a}$ Newton's second law (valid for both solids and fluids)

Weight = force on object due to gravity

$$W = mg$$
 $g = 9.81 \text{ m/s}^2$
= 32.2 ft/s²

SI:
$$W^{\circ}(N) = m \text{ (kg)} \cdot 9.81 \text{ m/s}^2$$

BG:
$$\mathcal{W}^{\circ}(lbf) = m(slug) \cdot 32.2ft/s^2$$

EE:
$$\mathcal{W}(lbf) = \frac{m(lbm)}{g_c} \cdot 32.2 \text{ ft/s}^2$$

$$g_c = 32.2 \frac{lbm \cdot ft}{s^2 \cdot lbf} = 32.2 \frac{lbm}{slug}, \text{ i.e., } 1 \text{ slug} = 32.2 \text{ lbm}$$

$$1 N = 1 kg \cdot 1 m/s^2$$

$$1 lbf = 1 slug \cdot 1 ft/s^2$$

System; Extensive and Intensive Properties

System = fixed amount of matter = mass m

Therefore, by definition

$$\frac{d(m)}{dt} = 0$$

Properties are further distinguished as being either extensive or intensive.

Extensive properties: depend on total mass of system,

e.g., m and W

Intensive properties: independent of amount of mass of

system, e.g., p (force/area) and ρ

(mass/volume)

Properties Involving the Mass or Weight of the Fluid

Specific Weight,
$$\gamma = \text{gravitational force (i.e., weight) per unit volume } \forall V = \mathcal{W}/V = mg/V = \rho g N/m^3$$

(Note that specific properties are extensive properties per unit mass or volume)

Mass Density
$$\rho = \text{mass per unit volume}$$

= $m/\forall \text{kg/m}^3$

Specific Gravity S = ratio of γ_{liquid} to γ_{water} at standard $T = 4^{\circ}C$ = $\gamma/\gamma_{water, 4^{\circ}C}$ dimensionless (or γ_{air} at standard conditions for gases)

 $\gamma_{\text{water, 4}^{\circ}\text{C}} = 9810 \text{ N/m}^3 \text{ for } T = 4^{\circ}\text{C} \text{ and atmospheric pressure}$

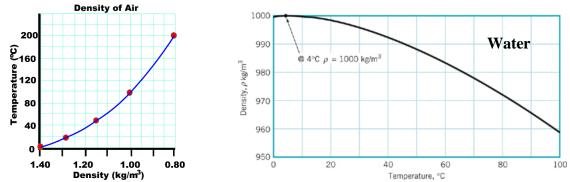
Variation in Density

gases:
$$\rho = \rho$$
 (gas, T , p) equation of state (p - v - T)
$$= p/RT \quad \text{ideal gas}$$

$$R = R \text{ (gas)}$$

e.g.
$$R$$
 (air) = 287.05 N·m/kg·°K
 ρ (air) = 1.225 kg/m³ at Standard Atmosphere
($T = 15$ °C and $p = 101.33$ kPa)

liquids: $\rho \sim \text{constant}$



Note: For a change of 100°C in temperature, density changes about 40% for air while about only 4% for water.

| Liquid and temperature | Density | Density |
|---------------------------------|------------|--------------------------|
| | (kg/m^3) | (slugs/ft ³) |
| Water 20°C (68°F) | 998 | 1.94 |
| Ethyl alcohol 20°C (68°F) | 799 | 1.55 |
| Glycerine 20°C (68°F) | 1,260 | 2.45 |
| Kerosene 20°C (68°F) | 814 | 1.58 |
| Mercury 20°C (68°F) | 13,350 | 26.3 |
| Sea water 10°C at 3.3% salinity | 1,026 | 1.99 |
| SAE 10W 38°C(100°F) | 870 | 1.69 |
| SAE 10W-30 8°C(100°F) | 880 | 1.71 |
| SAE 30 38°C(100°F) | 880 | 1.71 |

For greater accuracy can also use p-v-T diagram

$$\rho = \rho$$
 (liquid, T , p)

$$T \uparrow \qquad \rho \downarrow \\ p \uparrow \qquad \rho \uparrow$$

Properties Involving the Flow of Heat

For flows involving heat transfer such as gas dynamics additional thermodynamic properties are important, e.g.

specific heats c_p and c_v J/kg·°K

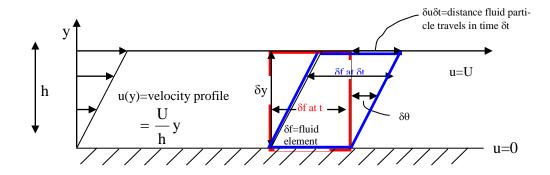
specific internal energy u J/kg

specific enthalpy $h = u + p/\rho$ J/kg

Viscosity

Recall definition of a fluid (substance that deforms continuously when subjected to a shear stress) and Newtonian fluid shear / rate-of-strain relationship: $\tau = \mu \dot{\theta}$.

Reconsider flow between fixed and moving parallel plates (Couette flow)



Newtonian fluid:
$$\tau = \mu \dot{\theta} = \mu \frac{\delta \theta}{\delta t}$$

 $\tan \delta \theta = \frac{\delta u \delta t}{\delta v}$ or $\delta \theta = \frac{\delta u \delta t}{\delta v}$ for small $\delta \theta$

therefore
$$\delta \dot{\theta} = \frac{\delta \mathbf{u}}{\delta \mathbf{y}}$$
 i.e., $\dot{\theta} = \frac{d\mathbf{u}}{d\mathbf{y}} = \text{velocity gradient}$

and
$$\tau = \mu \frac{du}{dy}$$

Exact solution for Couette flow is a linear velocity profile

$$u(y) = \frac{U}{h}y$$
 Note: $u(0) = 0$ and $u(h) = U$
i.e., satisfies no-slip
boundary condition

where

U/h = velocity gradient = rate of strain

 μ = coefficient of viscosity = proportionality constant for Newtonian fluid

$$\mu = \frac{\tau}{\frac{du}{dy}} = \frac{N/m^2}{\frac{m}{s}/m} = \frac{Ns}{m^2}$$

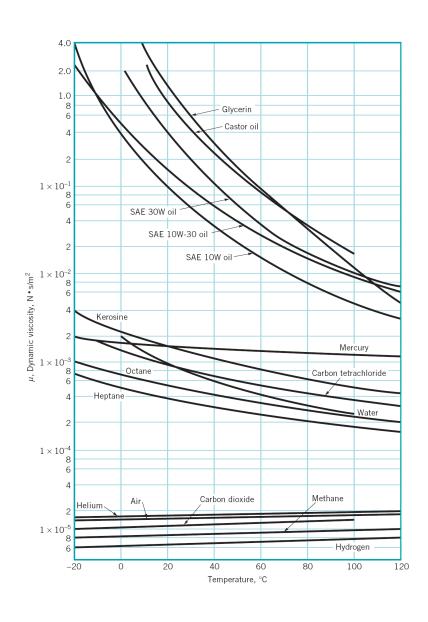
$$v = \frac{\mu}{\rho} = \frac{m^2}{s} = \text{kinematic viscosity}$$

 $\mu = \mu(\text{fluid}; T, p) = \mu(\text{gas/liquid}; T)$

gas and liquid $\mu \uparrow p \uparrow$, but small $\Delta \mu$

gas: $\mu \uparrow T \uparrow$ liquid: $\mu \downarrow T \uparrow$

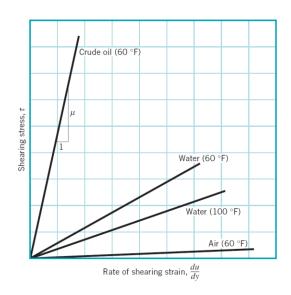
Due to structural differences, more molecular activity, decreased cohesive forces for gases

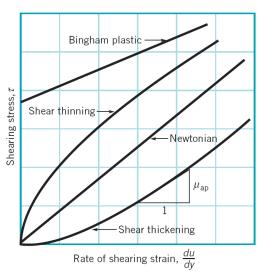


Newtonian vs. Non-Newtonian Fluids

Dilatant: $\tau \wedge du/dy \wedge$

Newtonian: $\tau \propto du/dy$ Pseudo plastic: $\tau \checkmark du/dy \uparrow$





$$\tau \propto \frac{du}{dy}$$

 μ = slope

$$\tau \propto \left(\frac{du}{dy}\right)^n$$

n > 1 slope increases with increasing τ (shear thickening)

n < 1 slope decreases with increasing τ (shear thinning) Ex) blood, paint, liquid plastic

Elasticity (i.e., compressibility)

Increasing/decreasing pressure corresponds to contraction/expansion of a fluid. The amount of deformation is called elasticity.

$$dp = -E_{v} \frac{dV}{V} \qquad dp > 0 \Longrightarrow \frac{dV}{V} < 0$$

Increase pressure, decrease volume. : minus sign used

$$E_V = -\frac{dp}{dV/V} = \frac{dp}{d\rho/\rho} = \rho \frac{dp}{d\rho} = \frac{N}{m^2}$$

Alternate form: $m = \rho \forall$ $dm = \rho d \forall + \forall d \rho = 0$ (by definition)

$$-\frac{d\Psi}{\Psi} = \frac{d\rho}{\rho}$$

Liquids are in general incompressible, e.g.

$$E_v = 2.2 \text{ GN/m}^2$$
 water
i.e. $\Delta \Psi = 0.05\%$ for $\Delta p = 1\text{MN/m}^2$
 $(\text{G=Giga=}10^9 \text{ M=Mega=}10^6 \text{ k=kilo=}10^3)$

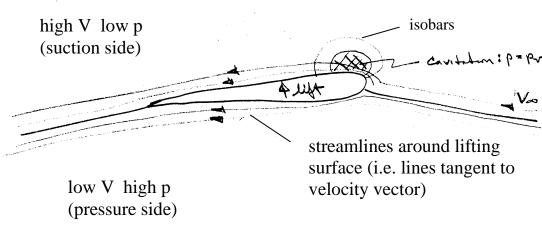
Gases are in general compressible, e.g. for ideal gas at T = constant (isothermal)

$$\frac{dp}{d\rho} = RT$$

$$E_{v} = \rho RT = p$$

Vapor Pressure and Cavitation

When the pressure of a liquid falls below the vapor pressure it evaporates, i.e., changes to a gas. If the pressure drop is due to temperature effects alone, the process is called boiling. If the pressure drop is due to fluid velocity, the process is called cavitation. Cavitation is common in regions of high velocity, i.e., low p such as on turbine blades and marine propellers.

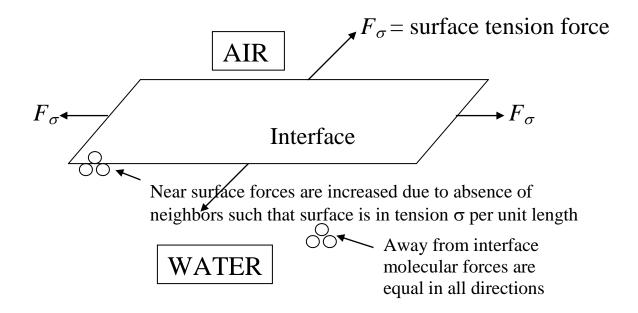


Cavitation number =
$$\frac{p - p_{v}}{\frac{1}{2}\rho V_{\infty}^{2}}$$

< 0 implies cavitation

Surface Tension and Capillary Effects

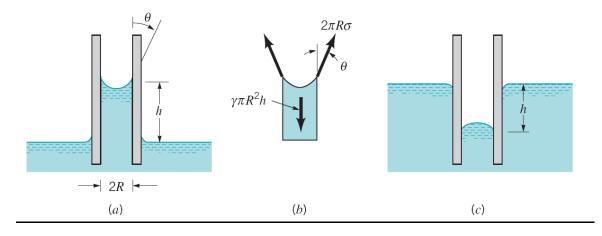
Two non-mixing fluids (e.g., a liquid and a gas) will form an interface. The molecules below the interface act on each other with forces equal in all directions, whereas the molecules near the surface act on each other with increased forces due to the absence of neighbors. That is, the interface acts like a stretched membrane



 $\sigma_{\text{air/water}} = 0.073 \text{ N/m}$ $F_{\sigma} = \sigma \times L = \text{line force with direction normal to the cut}$ L = length of cut through the interface

Effects of surface tension:

Contact angle:



$$\theta < 90^{\circ}$$
, Wetting e.g., Water, $\theta \approx 0^{\circ}$

 $\theta > 90^{\circ}$, Non-wetting e.g., Mercury, $\theta \approx 130^{\circ}$

1. Capillary action in small tube (See page 21)

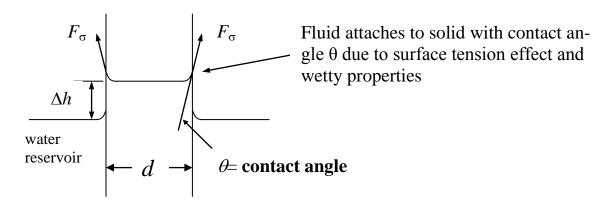
$$\Delta h = 4\sigma/\gamma d$$

2. Pressure difference across curved interface (See page 22)

$$\Delta p = \sigma/R$$
 R = radius of curvature

- 3. Transformation of liquid jet into droplets
- 4. Binding of wetted granular material such as sand
- 5. Capillary waves: surface tension acts as restoring force resulting in interfacial waves called capillary waves

Capillary tube



Example:

Capillary tube d = 1.6mm = 0.0016m $F_{\sigma} = \sigma \times L$, L=length of contact line between fluid & solid

water reservoir at 20° C, $\sigma = 0.073$ N/m, $\gamma = 9790$ N/m³

$$\Delta h = ?$$

$$\Sigma F_z = 0$$

$$F_{\sigma,z} - W^\circ = 0$$

$$\sigma \pi d \cos \theta - \rho g \Psi = 0 \qquad \theta \sim 0^\circ \Rightarrow \cos \theta = 1$$

$$\rho g = \gamma$$

$$\sigma \pi d - \gamma \Delta h \frac{\pi d^2}{4} = 0 \qquad \Psi = \Delta h \frac{\pi d^2}{4} \qquad \text{=Volume of fluid above reservoir}$$

$$\Delta h = \frac{4\sigma}{\gamma d} = 18.6 mm$$

Pressure jump across curved interfaces

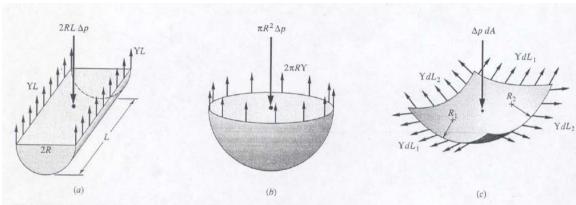


Fig. 1.9 Pressure change across a curved interface due to surface tension: (a) interior of a liquid cylinder; (b) interior of a spherical droplet; (c) general curved interface.

(a) Cylindrical interface Force Balance:

$$2\sigma L = 2 RL \cdot (p_i - p_o)$$

$$\Delta p = \sigma / R$$

 $p_i > p_o$, i.e. pressure is larger on concave vs. convex side of interface

(b) Spherical interface (Droplets)

$$\pi \cdot 2R\sigma = \pi R^2 \Delta p \quad \to \quad \Delta p = 2\sigma/R$$

(c) General interface

$$\Delta p = \sigma (R_1^{-1} + R_2^{-1})$$

$$R_{1,2} = \text{principal radii of curvature}$$

A brief history of fluid mechanics

See textbook section 1.10. (page 27)

Fluid Mechanics and Flow Classification

<u>Hydrodynamics:</u> flow of fluids for which density is constant such as liquids and low-speed gases. If in addition fluid properties are constant, temperature and heat transfer effects are uncoupled such that they can be treated separately.

Examples: hydraulics, low-speed aerodynamics, ship hydrodynamics, liquid and low-speed gas pipe systems

Gas Dynamics: flow of fluids for which density is variable such as high-speed gases. Temperature and heat transfer effects are coupled and must be treated concurrently.

Examples: high-speed aerodynamics, gas turbines, high-speed gas pipe systems, upper atmosphere